

Why Has There Been No Global Warming For The Past Decade?

Physics Department Colloquium
UNC Chapel Hill
September 8, 2014

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THE LONG-TERM IMPACTS OF INCREASING ATMOSPHERIC CARBON DIOXIDE LEVELS

Edited by

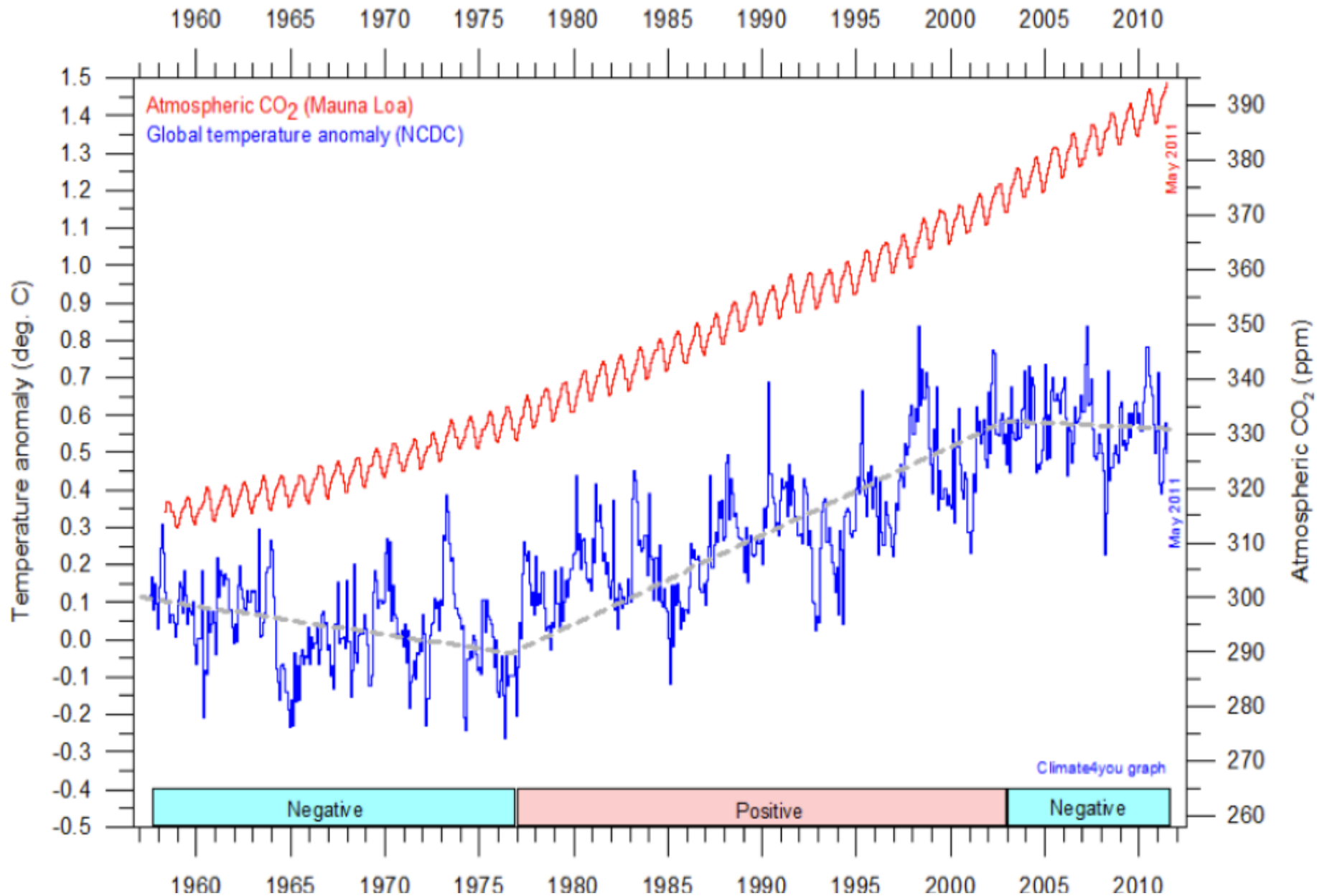
GORDON J. MACDONALD

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What is all the fuss about?



Climate Models Don't Work; Red is Observed; Bars are Models

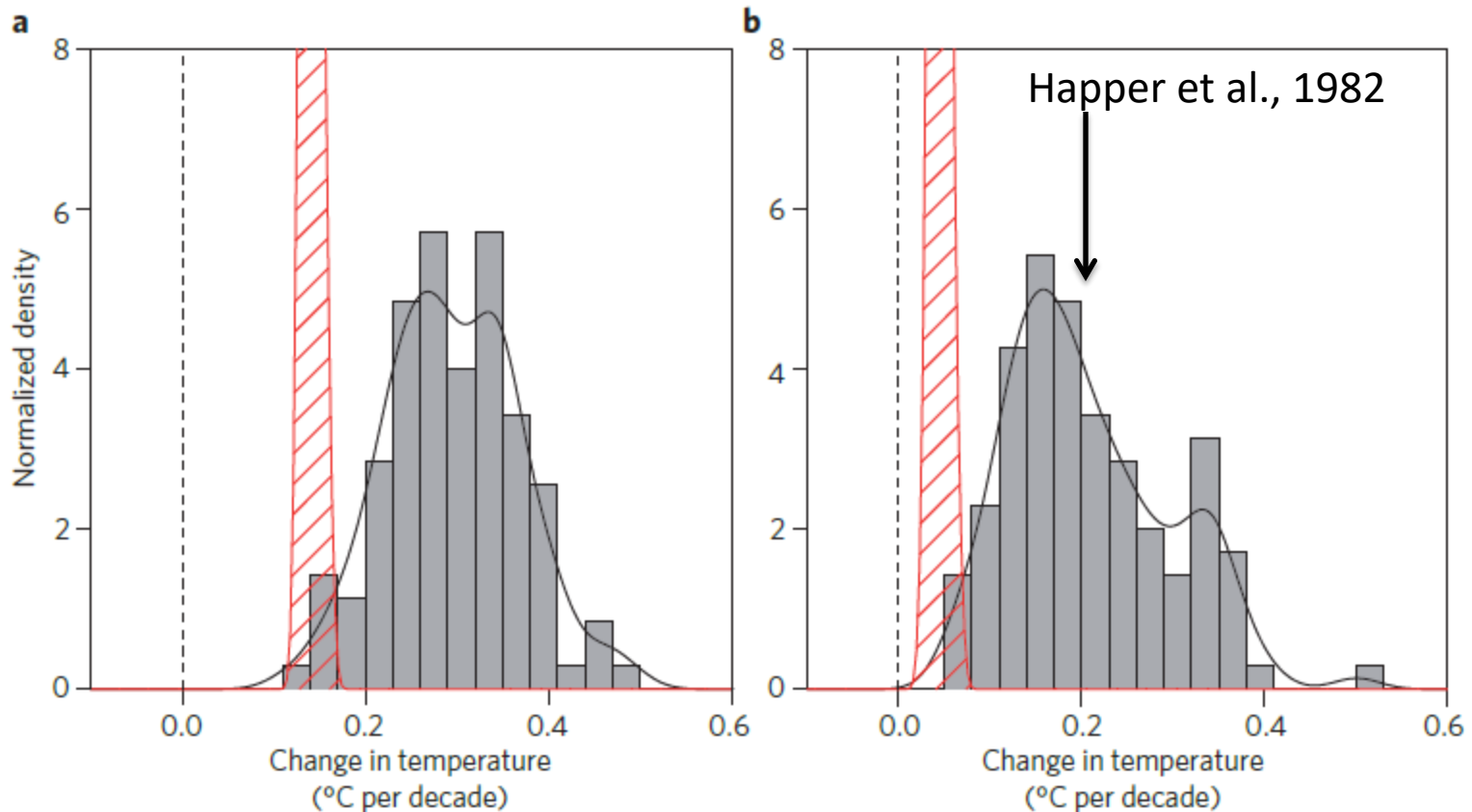
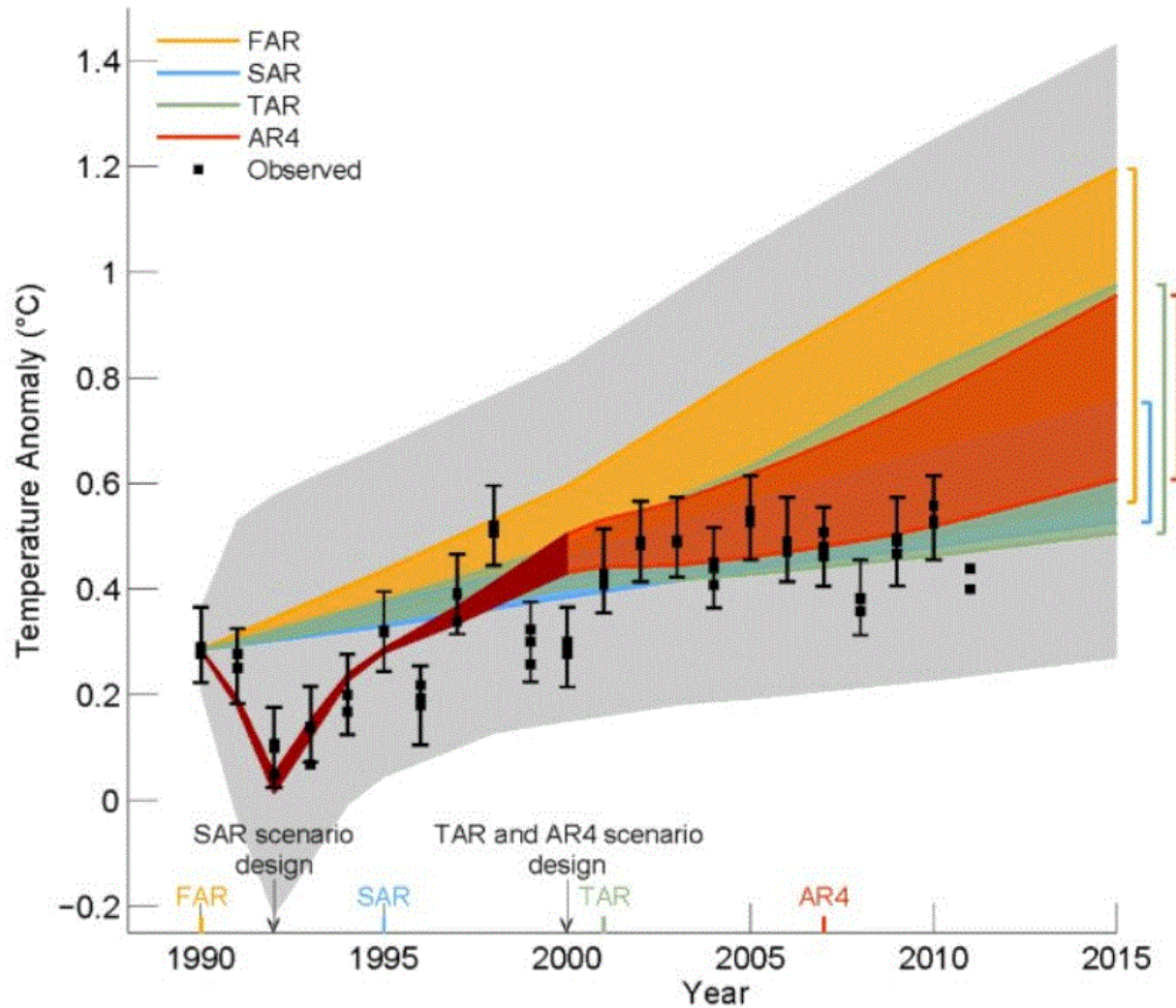


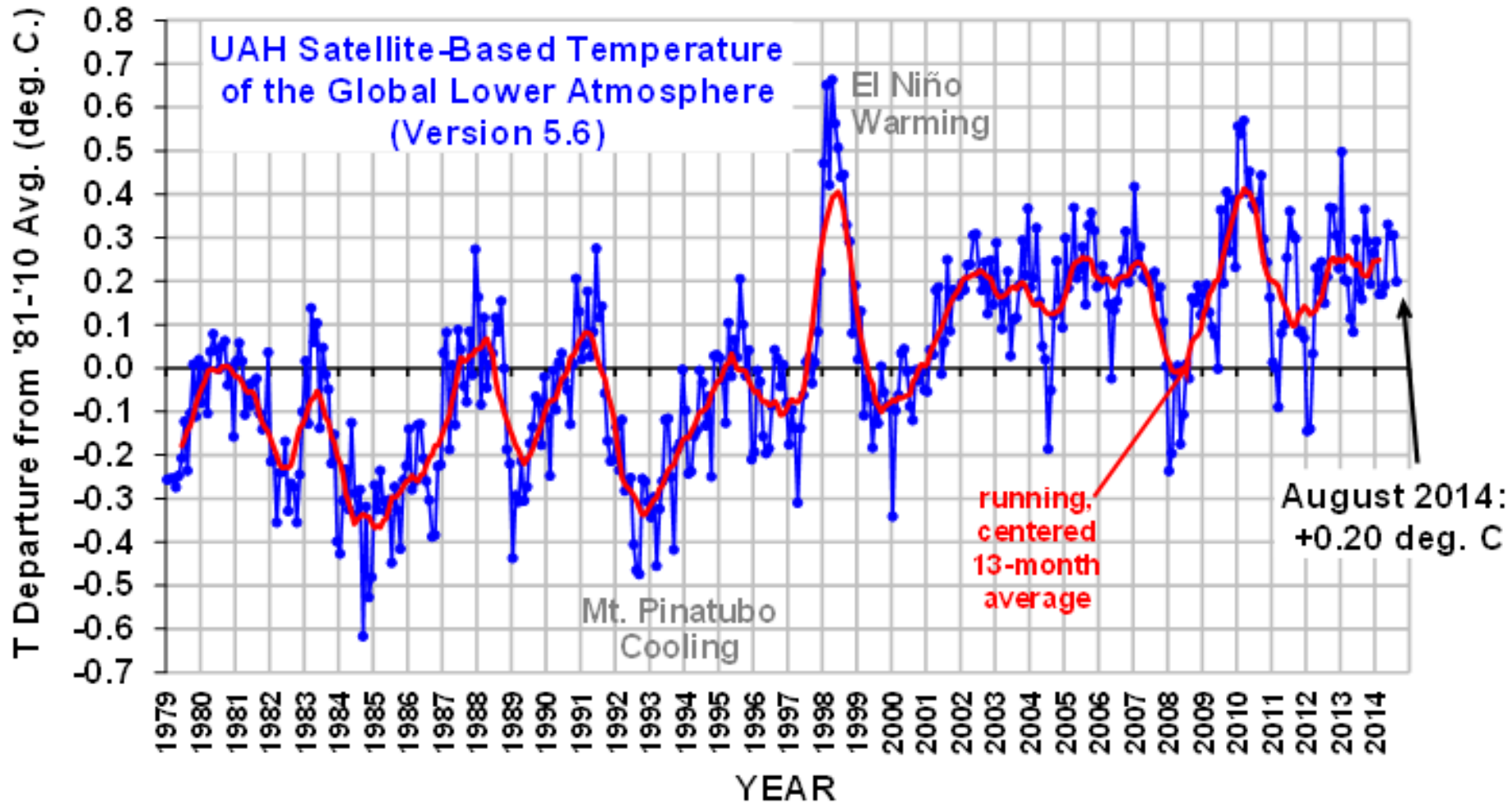
Figure 1 | Trends in global mean surface temperature. **a**, 1993–2012. **b**, 1998–2012. Histograms of observed trends (red hatching) are from 100 reconstructions of the HadCRUT4 dataset¹. Histograms of model trends (grey bars) are based on 117 simulations of the models, and black curves are smoothed versions of the model trends. The ranges of observed trends reflect observational uncertainty, whereas the ranges of model trends reflect forcing uncertainty, as well as differences in individual model responses to external forcings and uncertainty arising from internal climate variability.

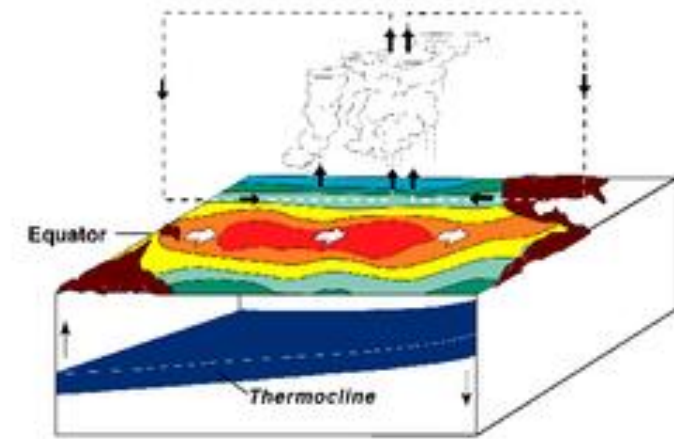
The press is beginning to notice!

[Klimawandel: Forscher rätseln über Stillstand bei Erderwärmung](#)

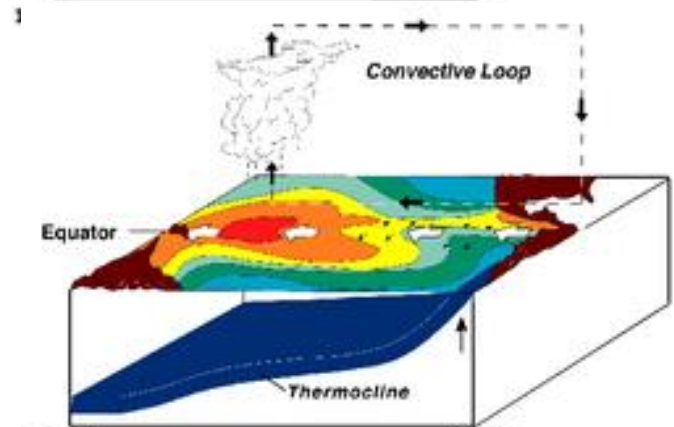


Satellite temperature measurements of lower atmosphere show little warming

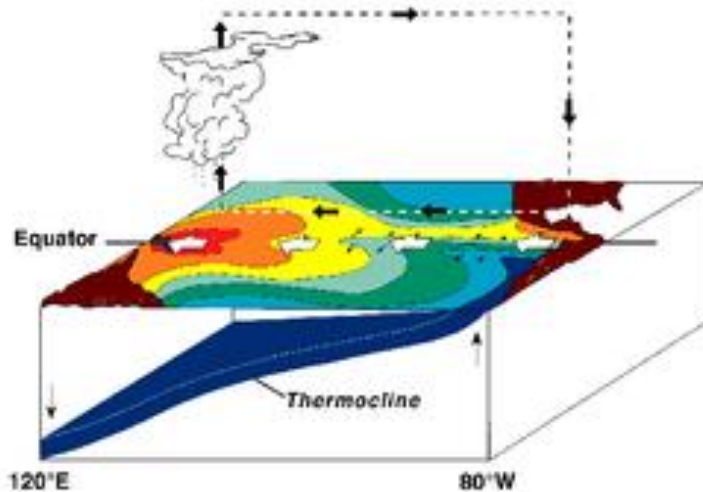




El Nino



Normal



La Nina

Al Gore, Climate Scientist, Nobel Laureate.

OUR CHOICE

A Plan to Solve the Climate Crisis

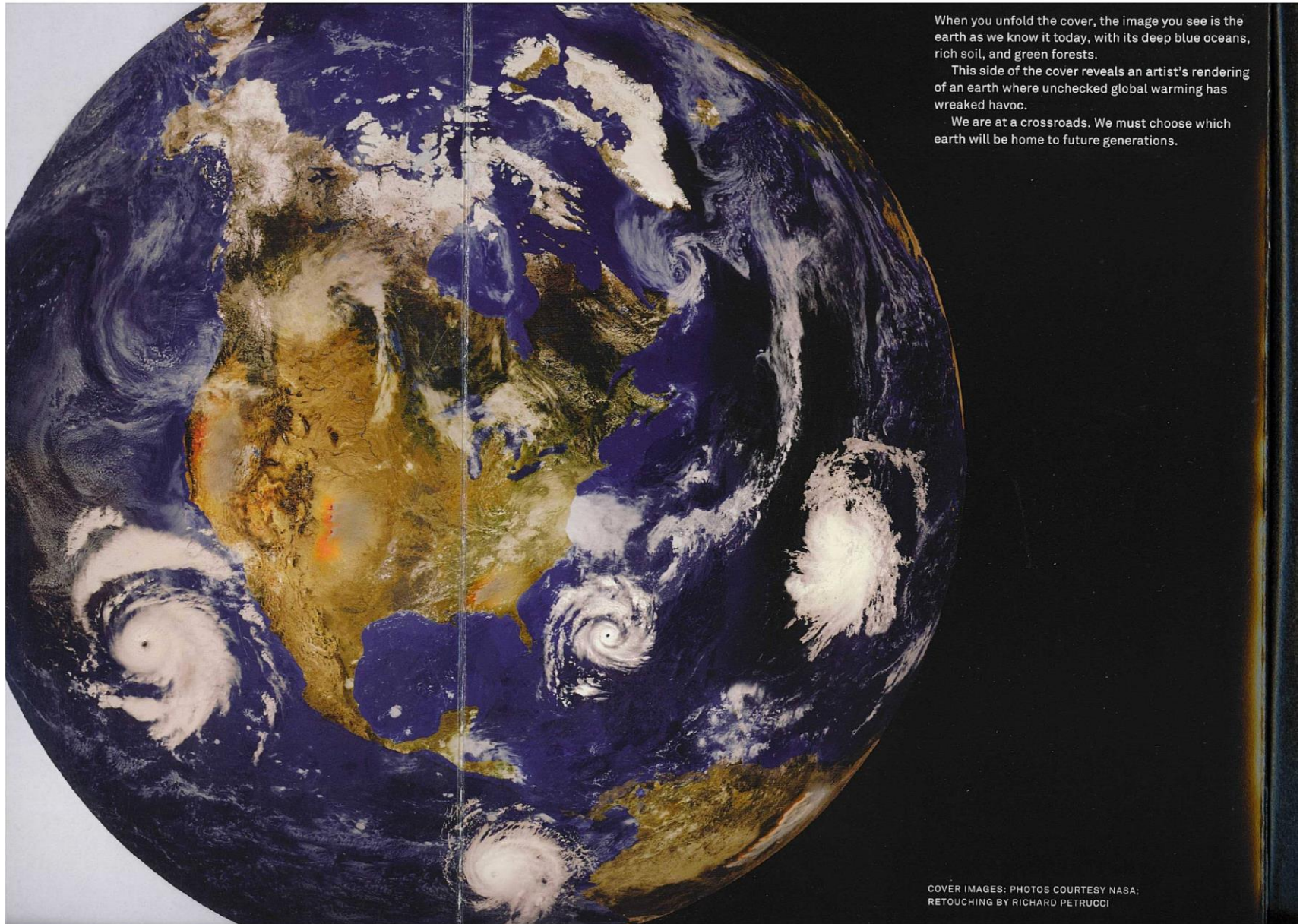


AL GORE

The Original for the Image on the Book's Cover.



The Image on the Inner Cover Showing the Lamentable Effects of More CO2

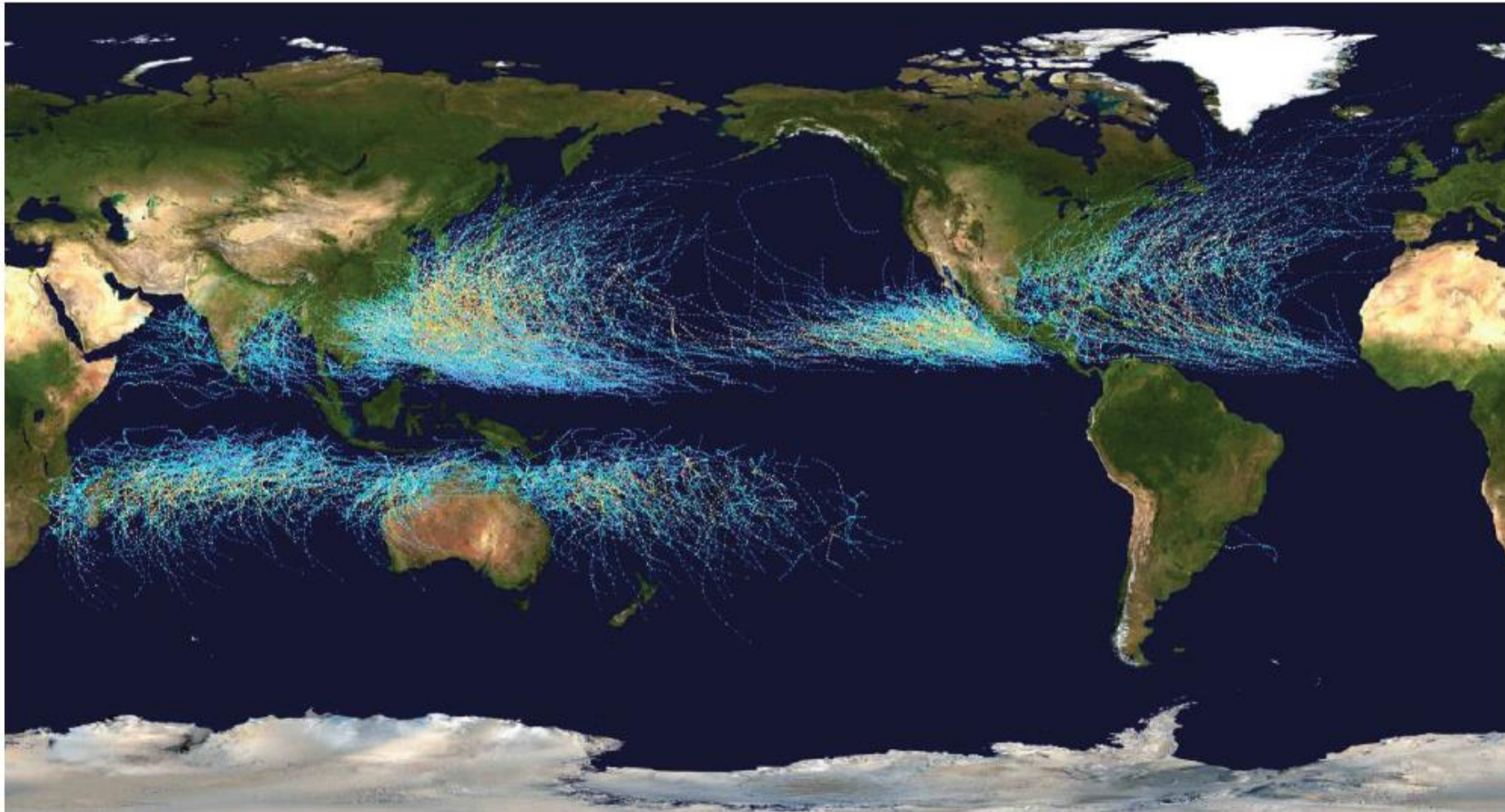


When you unfold the cover, the image you see is the earth as we know it today, with its deep blue oceans, rich soil, and green forests.

This side of the cover reveals an artist's rendering of an earth where unchecked global warming has wreaked havoc.

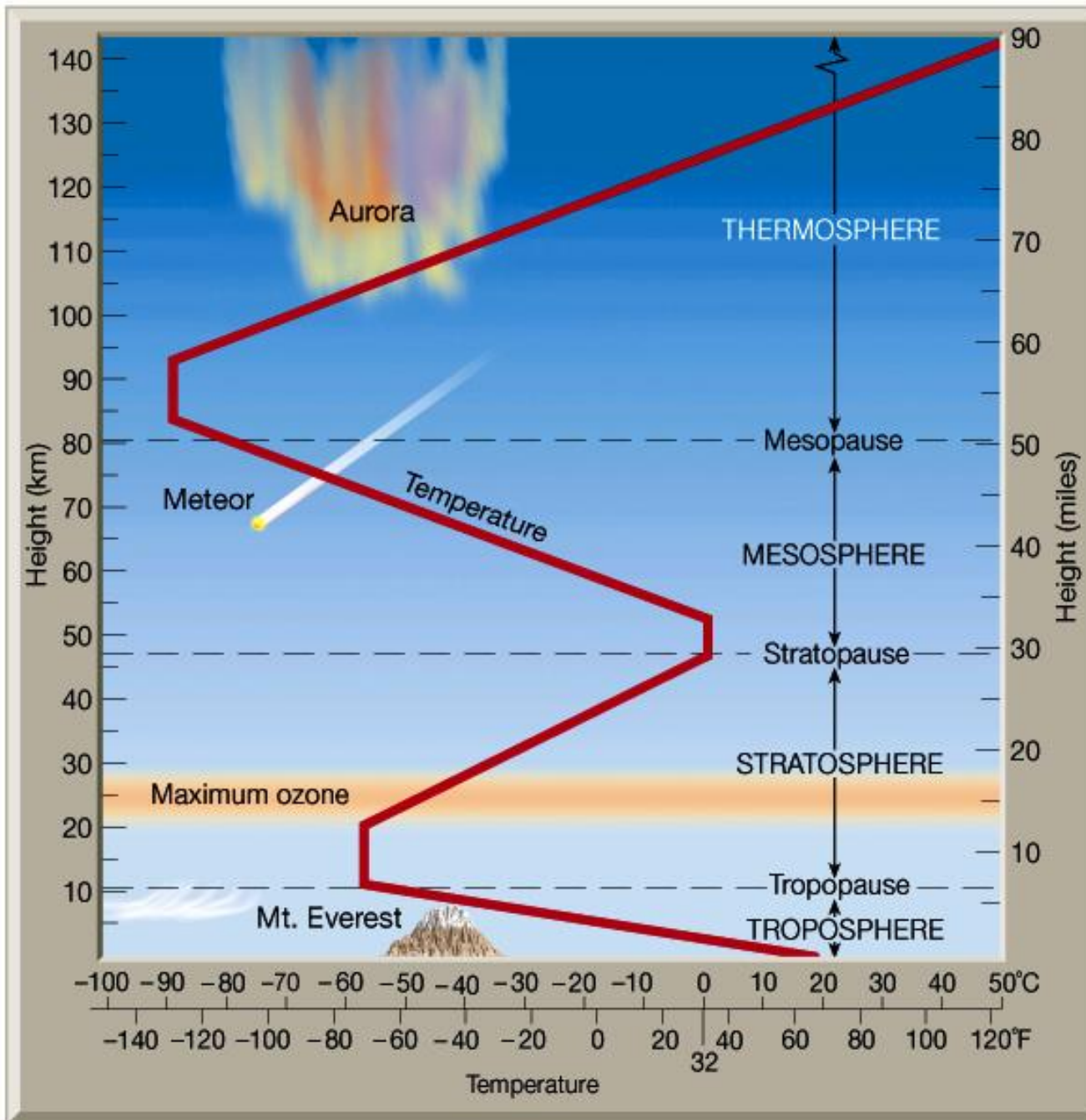
We are at a crossroads. We must choose which earth will be home to future generations.

Hurricane (Tropical Cyclone) Tracks; 1985-2005; (Wikipedia)

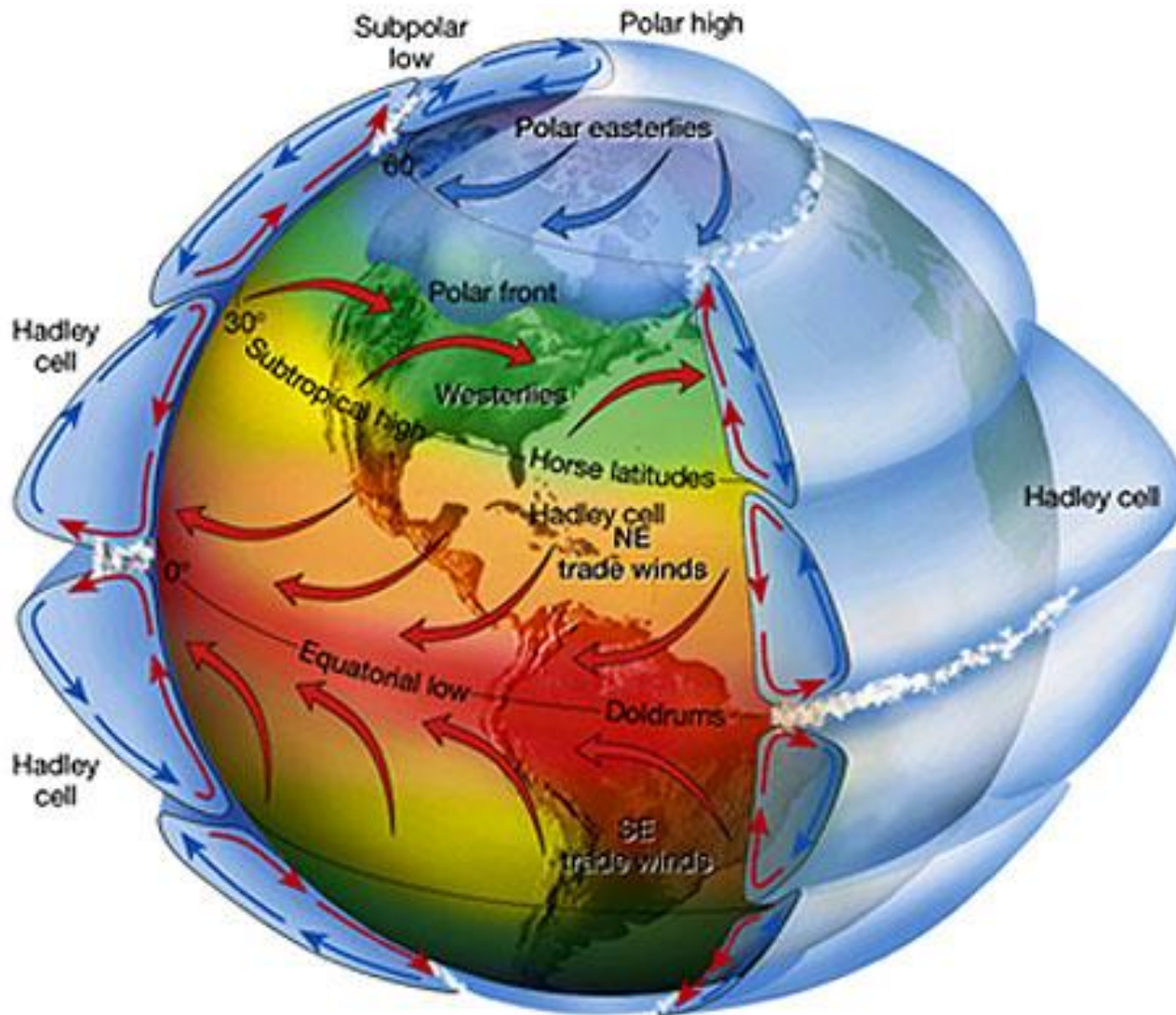


No Coriolis force at equator → No hurricanes

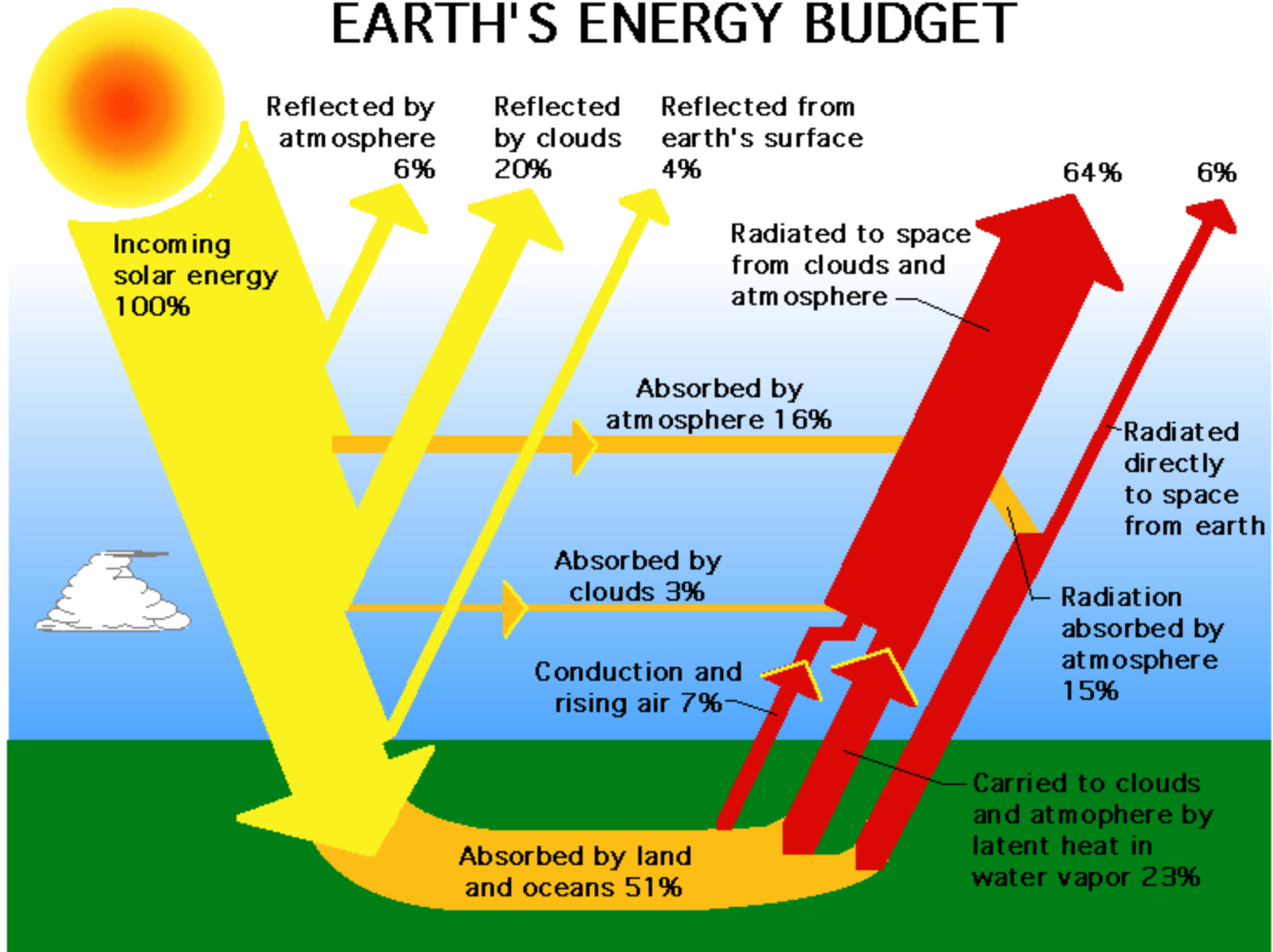
Low clouds spiral in counterclockwise in North; clockwise in South



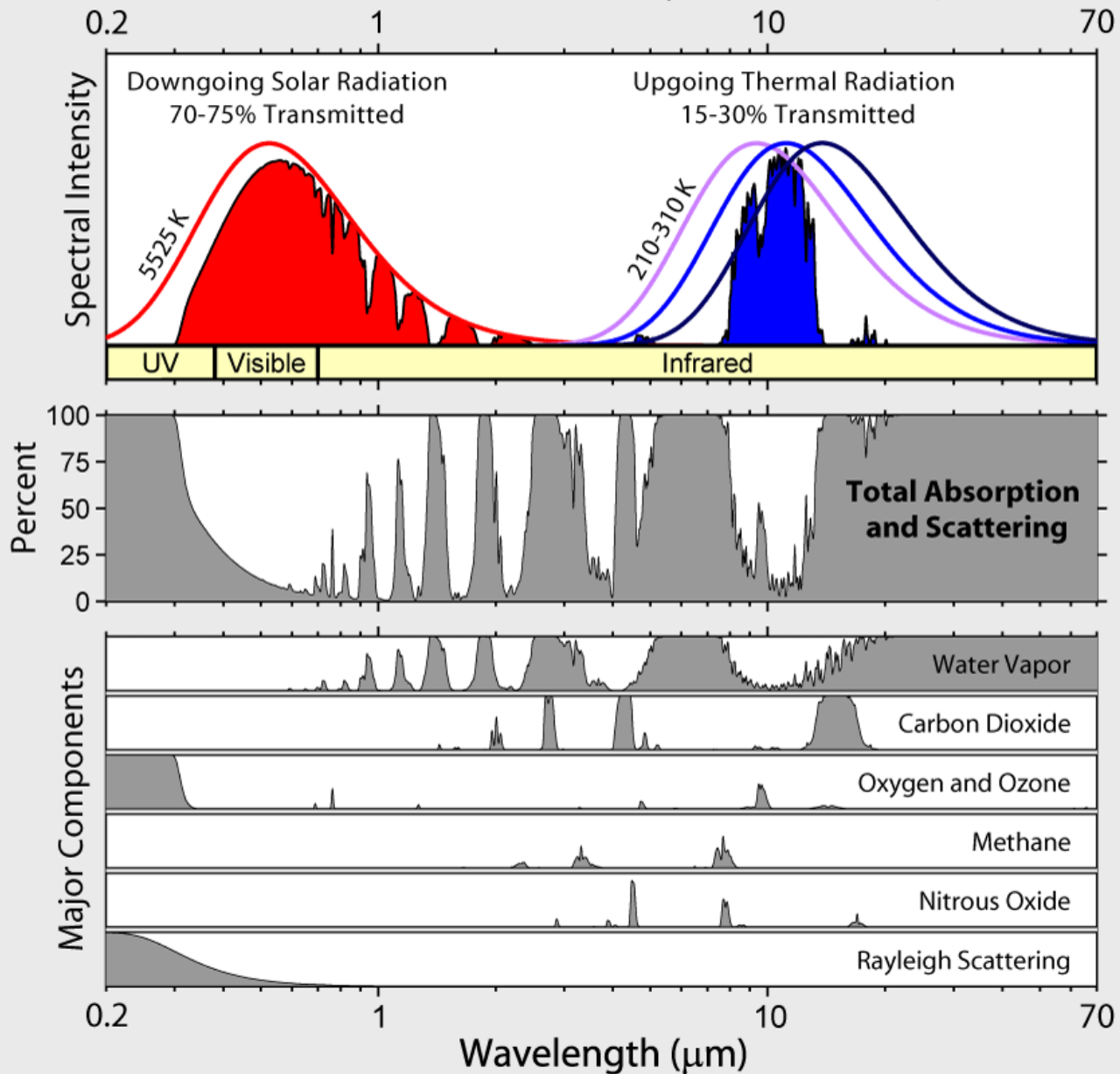
Atmospheric Circulation



EARTH'S ENERGY BUDGET



Radiation Transmitted by the Atmosphere



Examples of FTIR Data from a Satellite

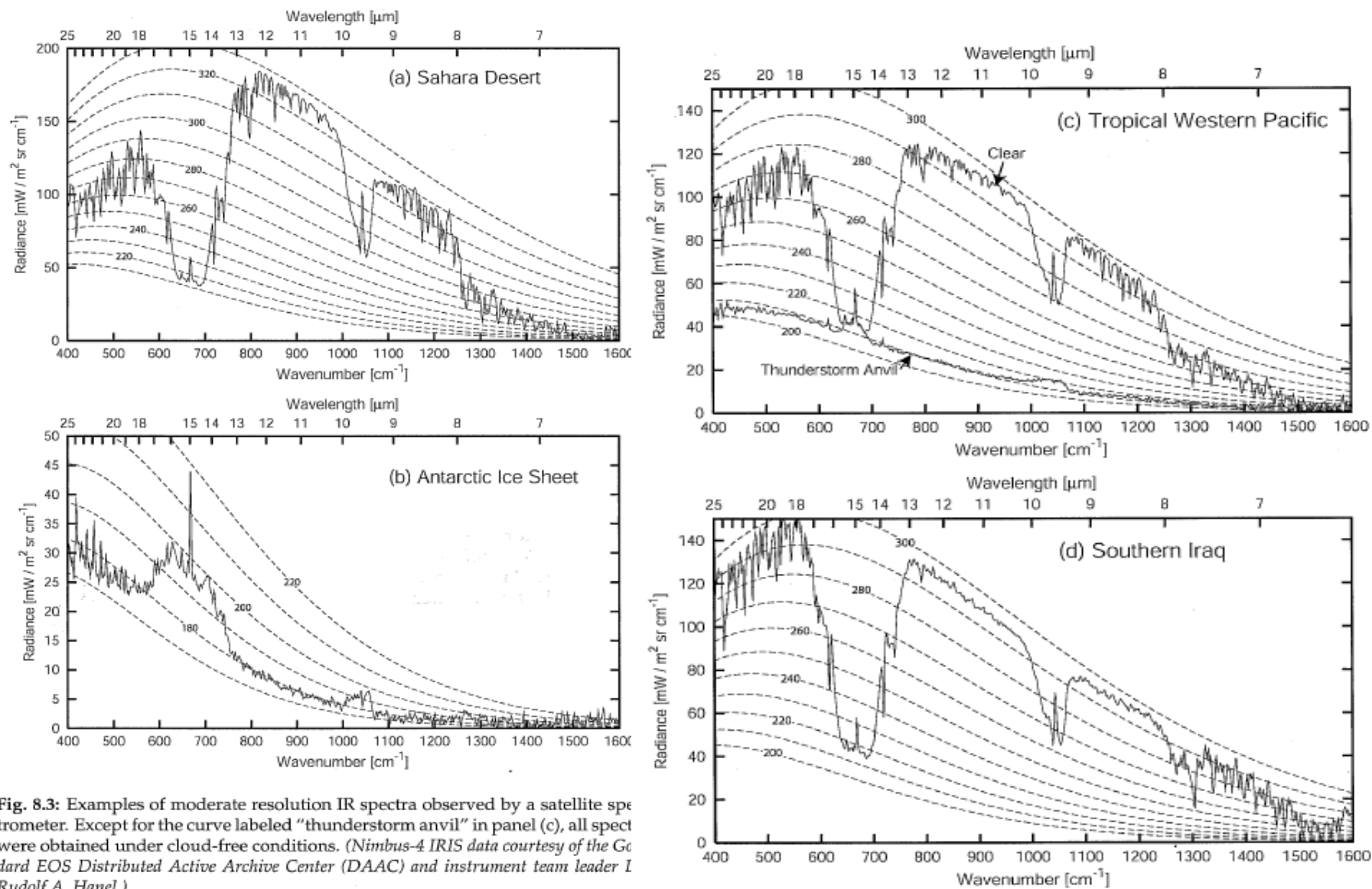


Fig. 8.3: Examples of moderate resolution IR spectra observed by a satellite spectrometer. Except for the curve labeled "thunderstorm anvil" in panel (c), all spectra were obtained under cloud-free conditions. (*Nimbus-4 IRIS data courtesy of the Goddard EOS Distributed Active Archive Center (DAAC) and instrument team leader I. Rudolf A. Hanel.*)

One Slide Summary of Global Warming Theory

Key Parameters: forcing ΔQ_2 and feedback f

Steady-State Temperature Change for Doubled CO2

$$\Delta T = \frac{T_e \Delta Q_2}{4Q_e(1-f)} = \frac{0.30 \Delta Q_2}{1-f} \text{ K m}^2 \text{ W}^{-1} = 3.4 \text{ K} = \frac{1.1 \text{ K}}{(1-2/3)}.$$

$$Q_e = (1 - \alpha_e)F_e/4 = \sigma \epsilon_e T_e^4 = 236 \text{ W m}^{-2} = \text{mean radiation of earth to space.}$$

$$\alpha_e = 0.306 \text{ albedo (mostly clouds).}$$

$$F_e = 1361 \text{ W m}^{-2} = \text{mean solar flux at earth's orbit.}$$

$$T_e = 288 \text{ K} = \text{mean surface temperature of Earth.}$$

$$\epsilon_e = 0.606 = \text{effective (contrived) emissivity of Earth for } T_e.$$

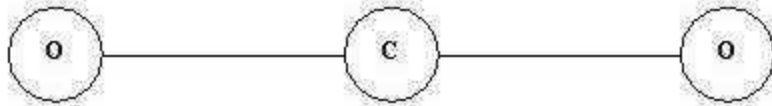
$$f = 0.25 T_e \partial \ln(1 - \alpha_e)/\partial T - T_e \partial \ln \epsilon_e/\partial T = 2/3 \text{ (per IPCC) = feedback.}$$

$$\Delta Q_2 = 3.7 \text{ W m}^{-2} \text{ (per IPCC) = radiative forcing for doubled CO2.}$$

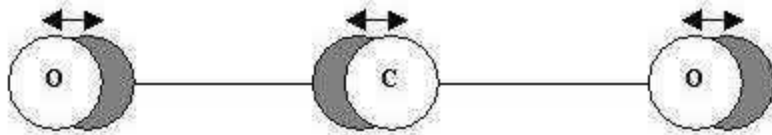
IPCC values of f and ΔQ_2 give far (at least a factor of 3) too much warming.

Both probably wrong in a way that exaggerates warming!

The Villain!



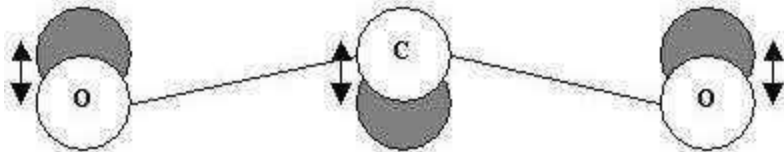
Molecular structure of Carbon Dioxide



The asymmetric stretch mode

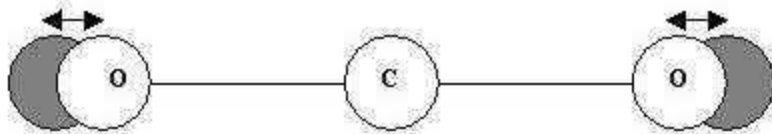
2349 cm^{-1}

frequency too high for greenhouse warming



The bending mode

666 cm^{-1} (Satan's number)
greenhouse warming

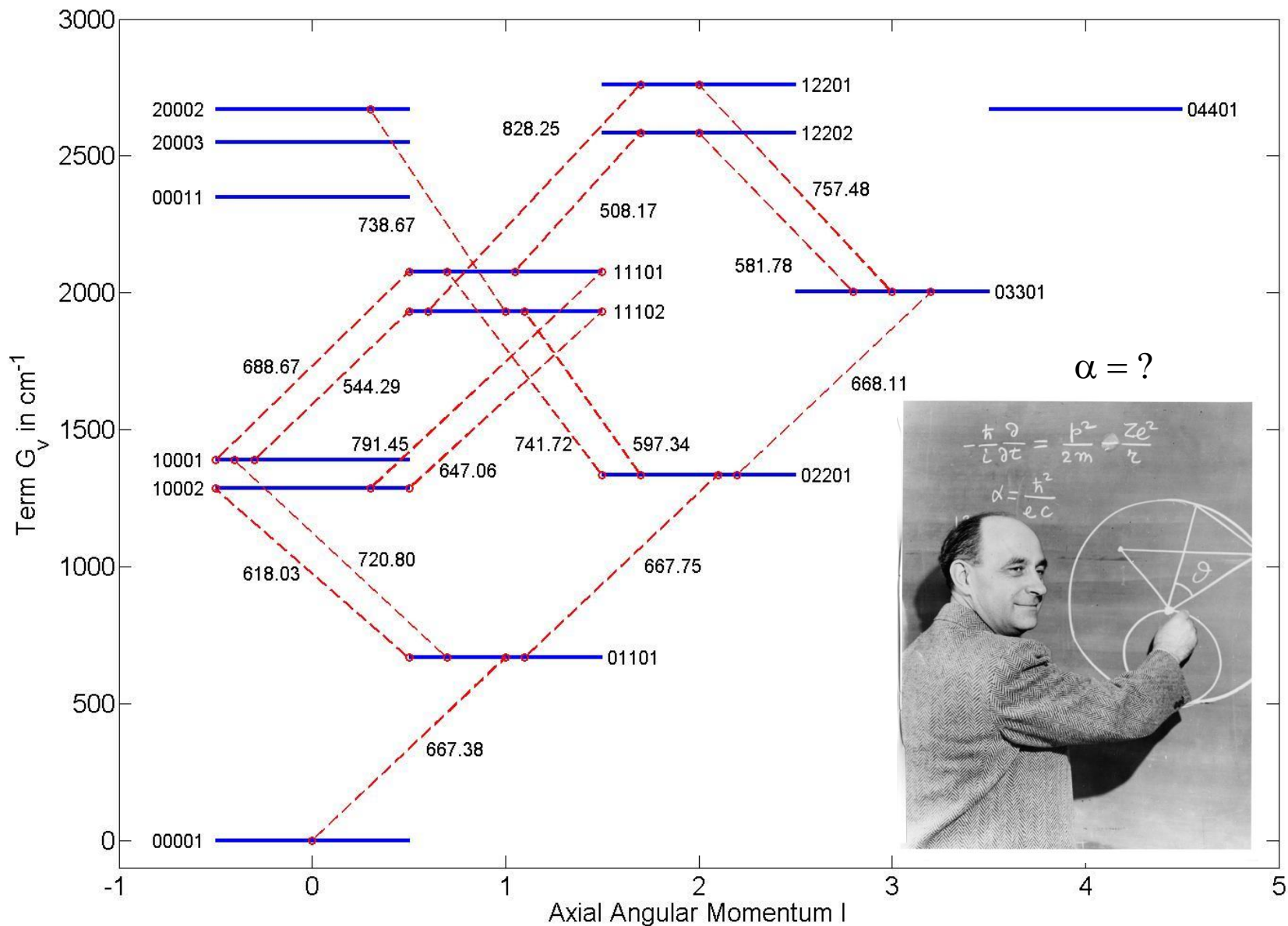


The symmetric stretch mode

1388 cm^{-1}

no changing dipole moment,
no absorption or emission

Fermi Resonances



What is wrong with this slide?

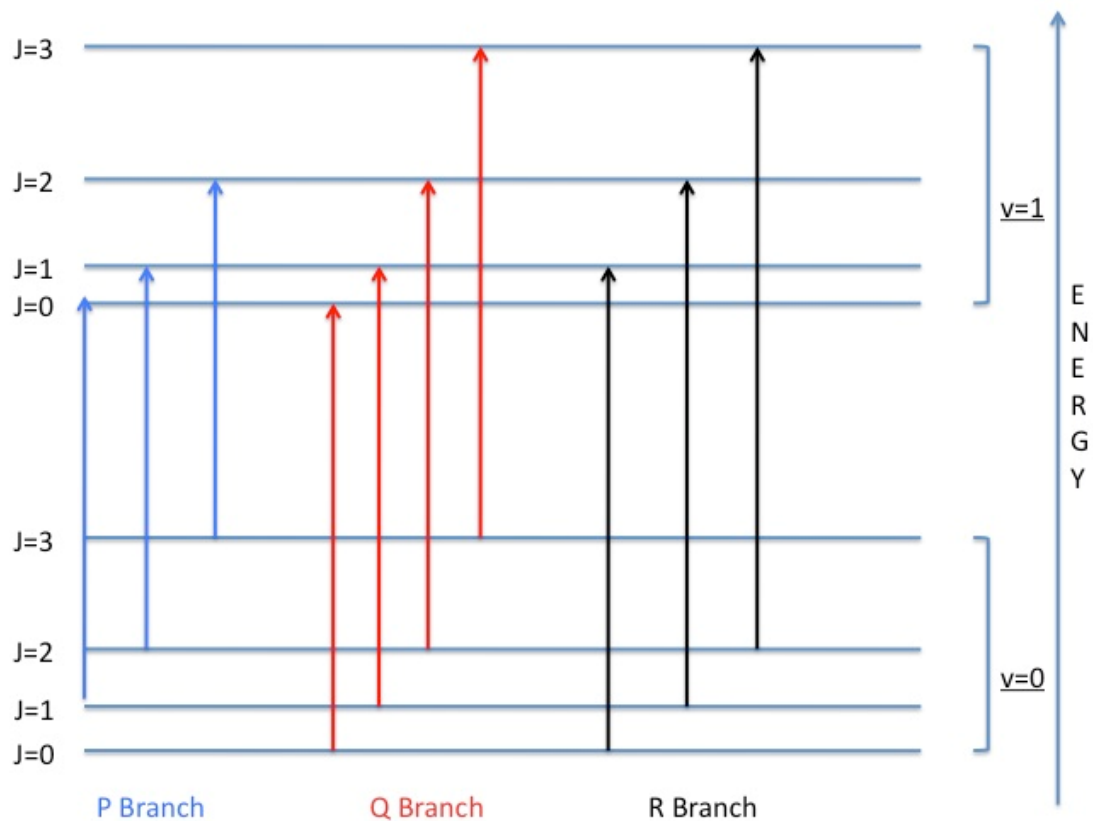


Fig. 6 Schematic diagram of P, Q, and R branch transitions

Attenuation Coefficient (e-foldings / length)

$$\kappa = N \sum_{eg} \sigma_{eg} = N \sum_{eg} S_{eg} G_{eg} \leftarrow \text{Line shape}$$

N is the number density of CO_2 molecules, and σ_{eg} is the cross section of the transition from a lower state g to an upper state e ; G_{eg} is the lineshape function.

Line strength (in cm)

$$S_{eg} = \frac{8\pi^3 \nu_{eg} |D_{eg}|^2 e^{-E_g/kT} (1 - e^{-hc\nu_{eg}/kT})}{hcZ}$$

ν_{eg} = frequency (in cm^{-1}) of the transition, D_{eg} = the electric dipole matrix element, E_g = lower-state energy, T = absolute temperature, k = Boltzmann's constant, h = Planck's constant, c = speed of light, Z = partition function.

Partition function

$$Z = \sum_j e^{-E_j/kT}$$

Normalized line shapes

$$\int_0^{\infty} G_{eg} d\nu = 1$$

A Lorentzian line shape

$$G_{eg} = \frac{\mu_{eg}/\pi}{\mu_{eg}^2 + (\nu - \nu_{eg})^2}$$

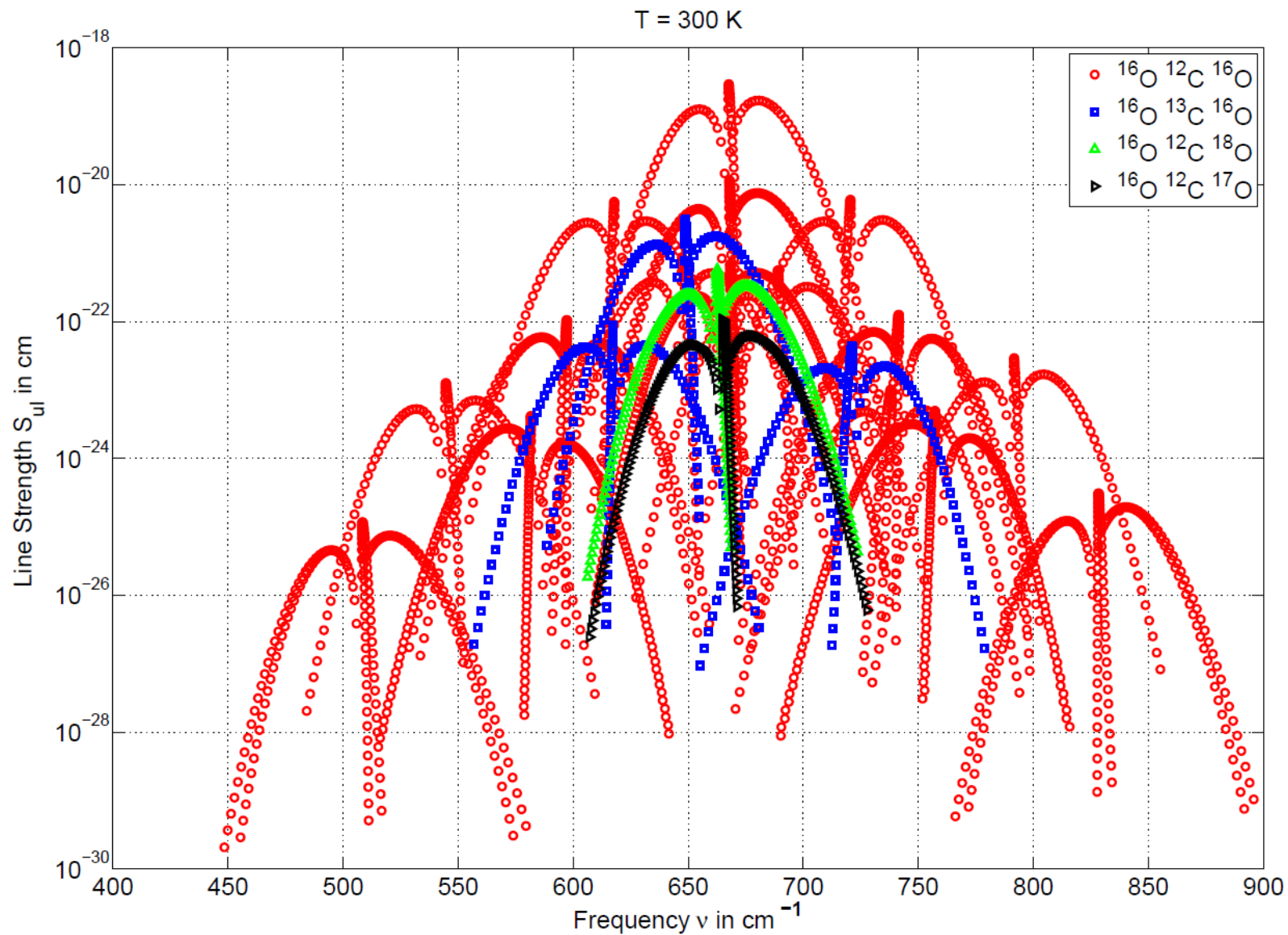
μ_{eg} = broadening; ν = frequency; ν_{eg} = resonance.

A Voigt line shape

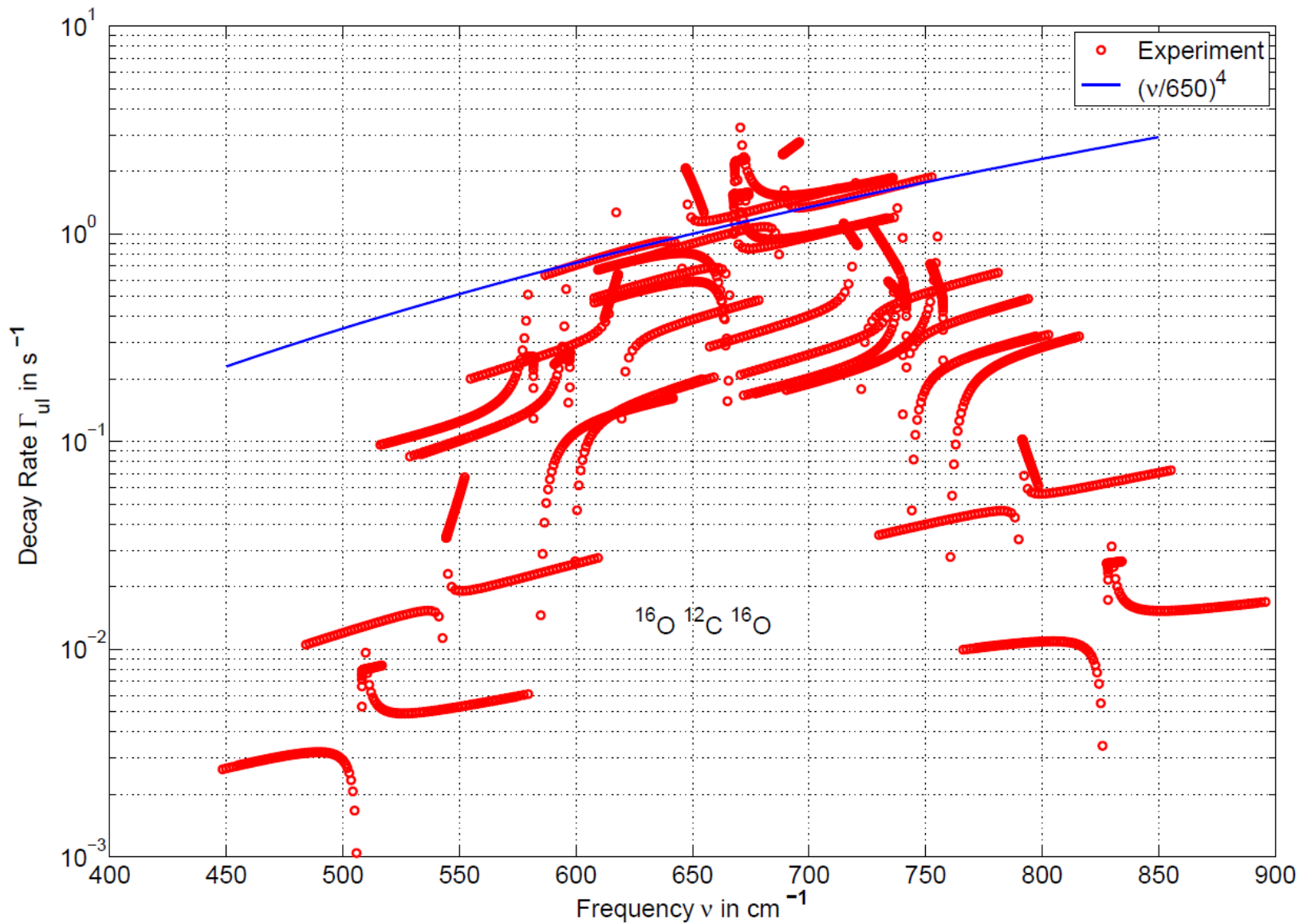
$$G_{eg} = \frac{\mu_{eg}}{\pi} \sqrt{\frac{m}{2\pi kT}} \int_{-\infty}^{\infty} \frac{e^{-mv^2/2kT} dv}{\mu_{eg}^2 + (\nu - \nu_{eg}[1 + v/c])^2}$$

Neither Lorentzian nor Voigt line shapes are correct in the far wings!

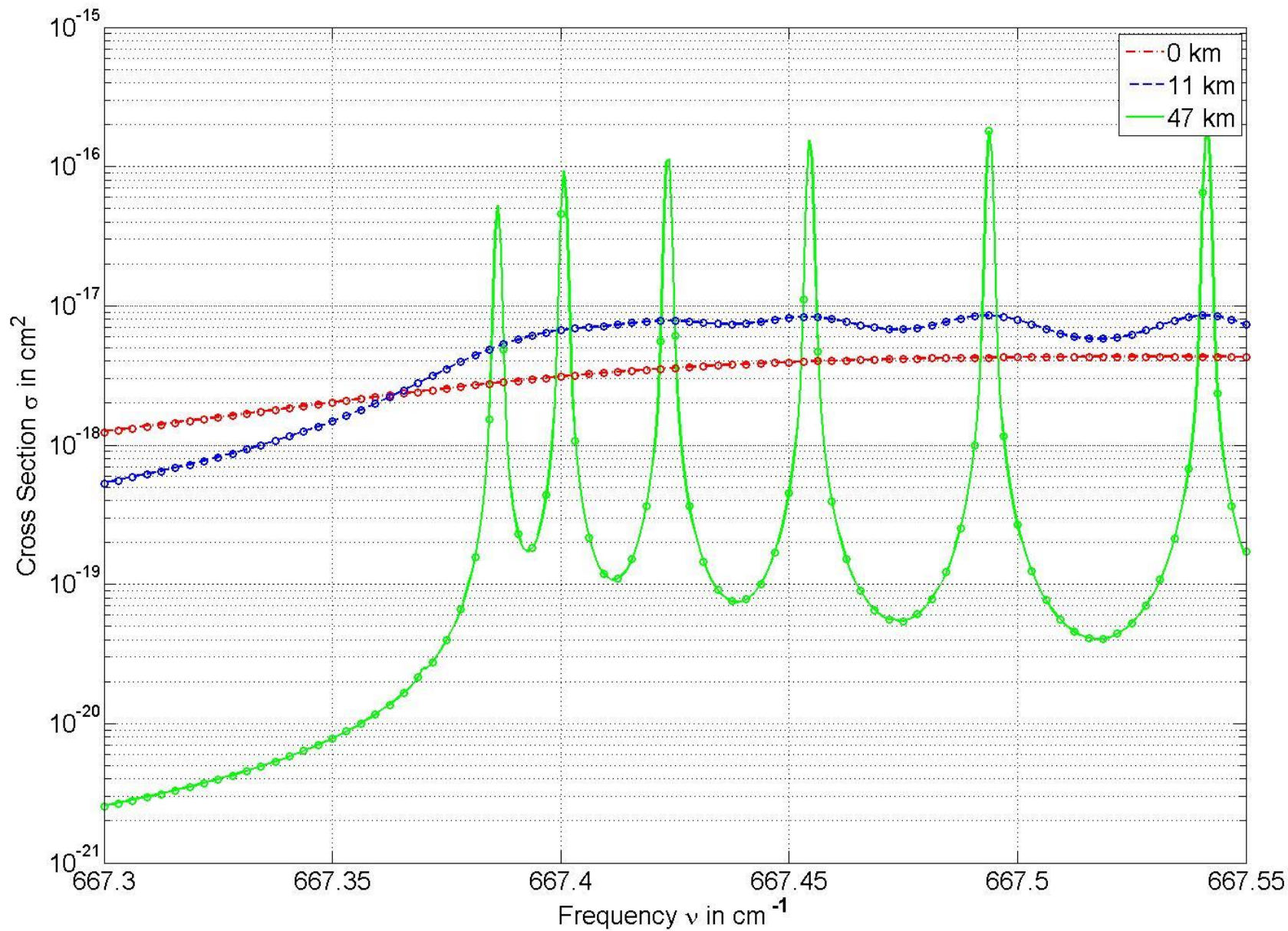
Thousands of lines! Linestrengths vary by 10 orders of magnitude.



Very Slow Partial Radiative Decay Rates of CO₂ Molecules



Less pressure broadening of cross sections at high altitude. Q-branch lines.



Schwartzschild-Milne Equation: Downwelling Flux at the Surface

$$J(0) = \int_0^{\infty} \kappa(z) B(z) e^{-\rho(z)} dz$$

Optical Depth from Surface
to Altitude z

$$\rho(z) = \int_0^z \kappa(z') dz'$$

Planck Brightness

$$B = \frac{2hc^2\nu^3}{e^x - 1}, \quad \text{where} \quad x = \frac{hc\nu}{kT}$$



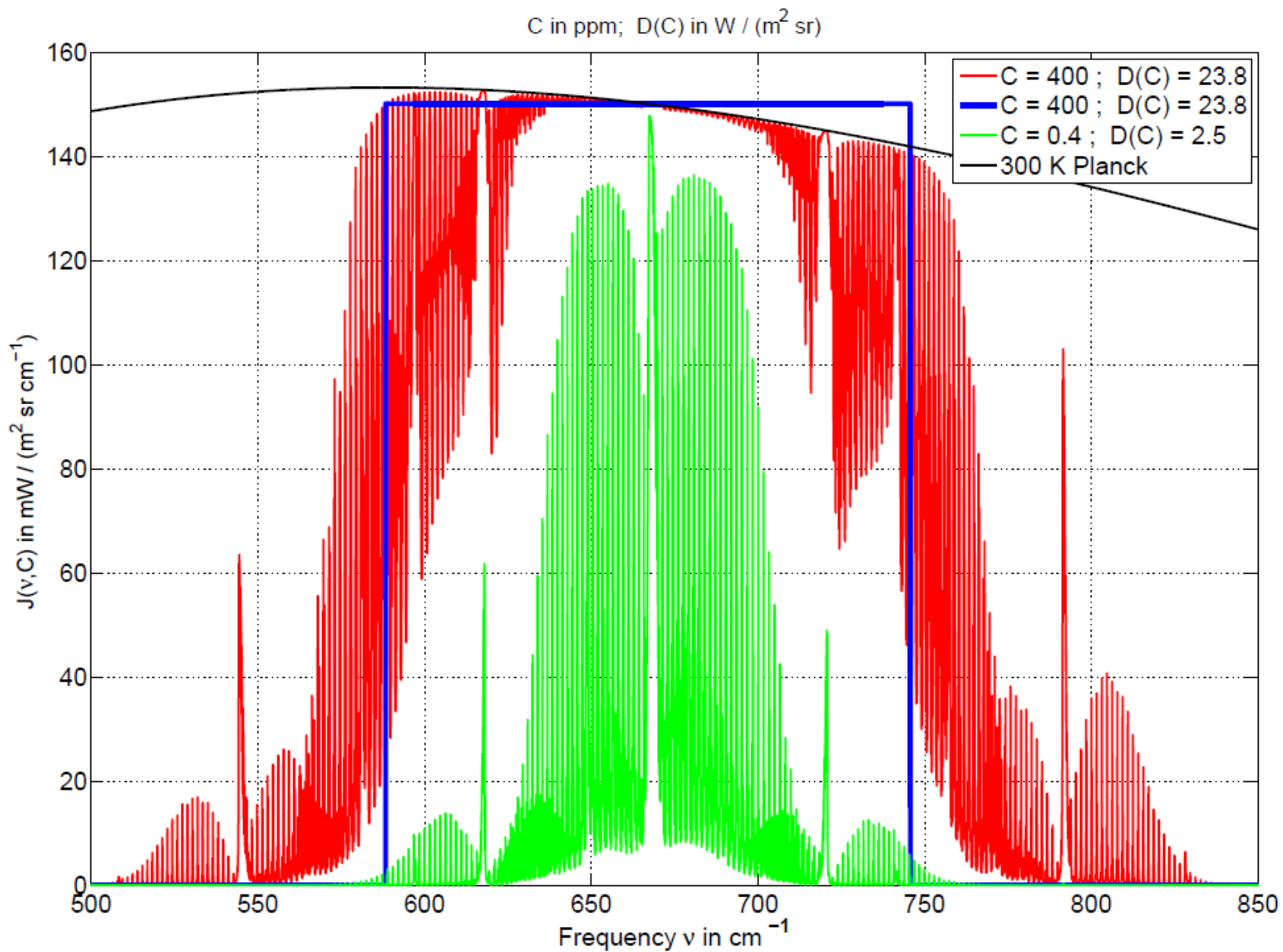
K. Schwartzschild



E. A. Milne

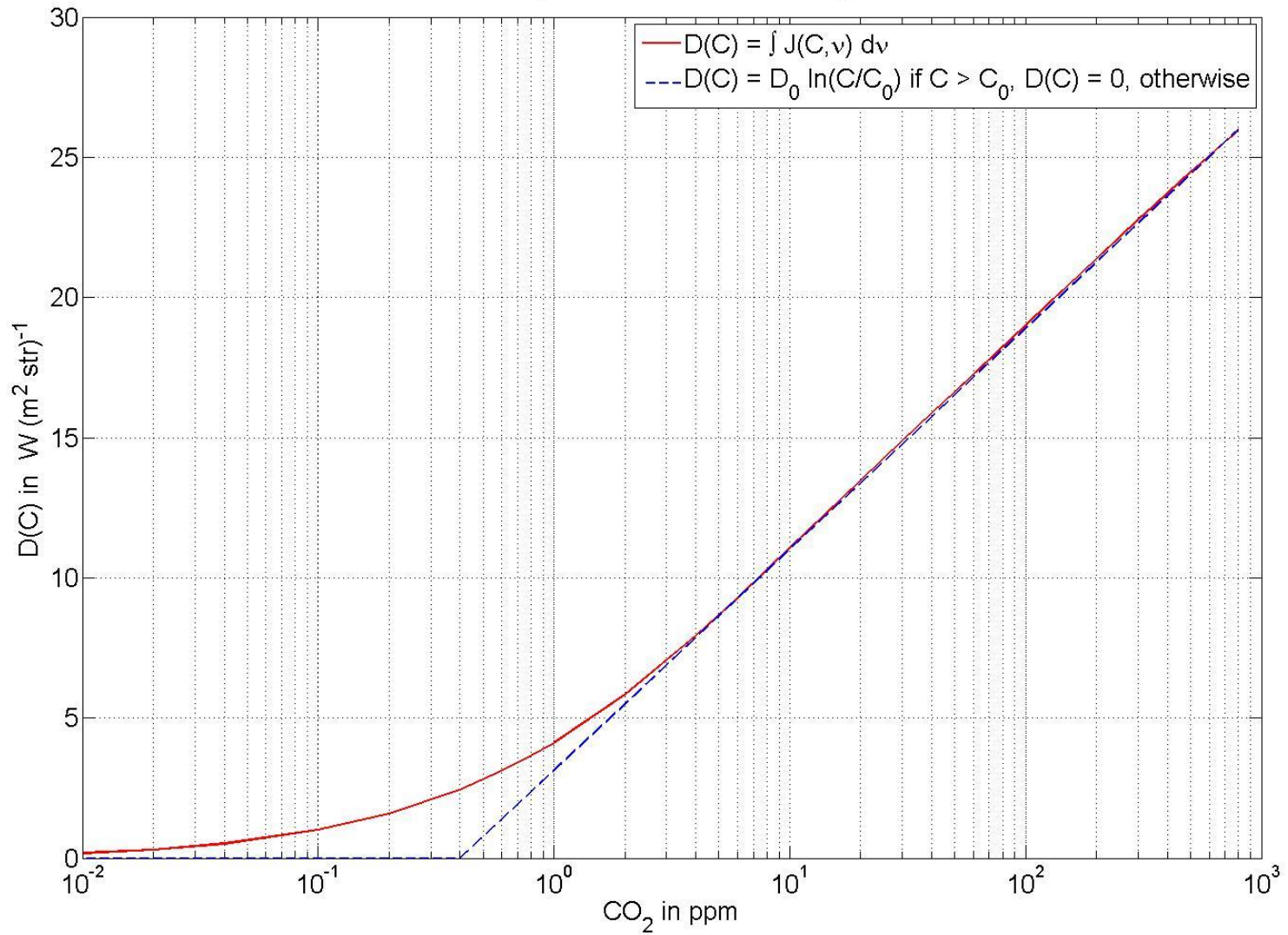


M. Planck



Logarithmic dependence of down-welling IR on CO₂ concentration.

$T_0 = 300 \text{ K}$; $D_0 = 3.42 \text{ W (m}^2 \text{ str)}^{-1}$ and $C_0 = 0.4 \text{ ppm}$



Logarithmic Dependence of Warming on CO₂ Concentration

Svante Arrhenius

Worlds in the Making, The Evolution of the Universe

(page 53)

“If the quantity of carbon dioxide in the air should sink to one half its present percentage, the temperature would fall by 4 K; a diminution by one-quarter would reduce the temperature by 8 K. On the other hand any doubling of the percentage of Carbon dioxide in the air would raise the temperature of the Earth’s surface by 4 K, and if the carbon dioxide were increased four fold, the temperature would rise by 8 K.”

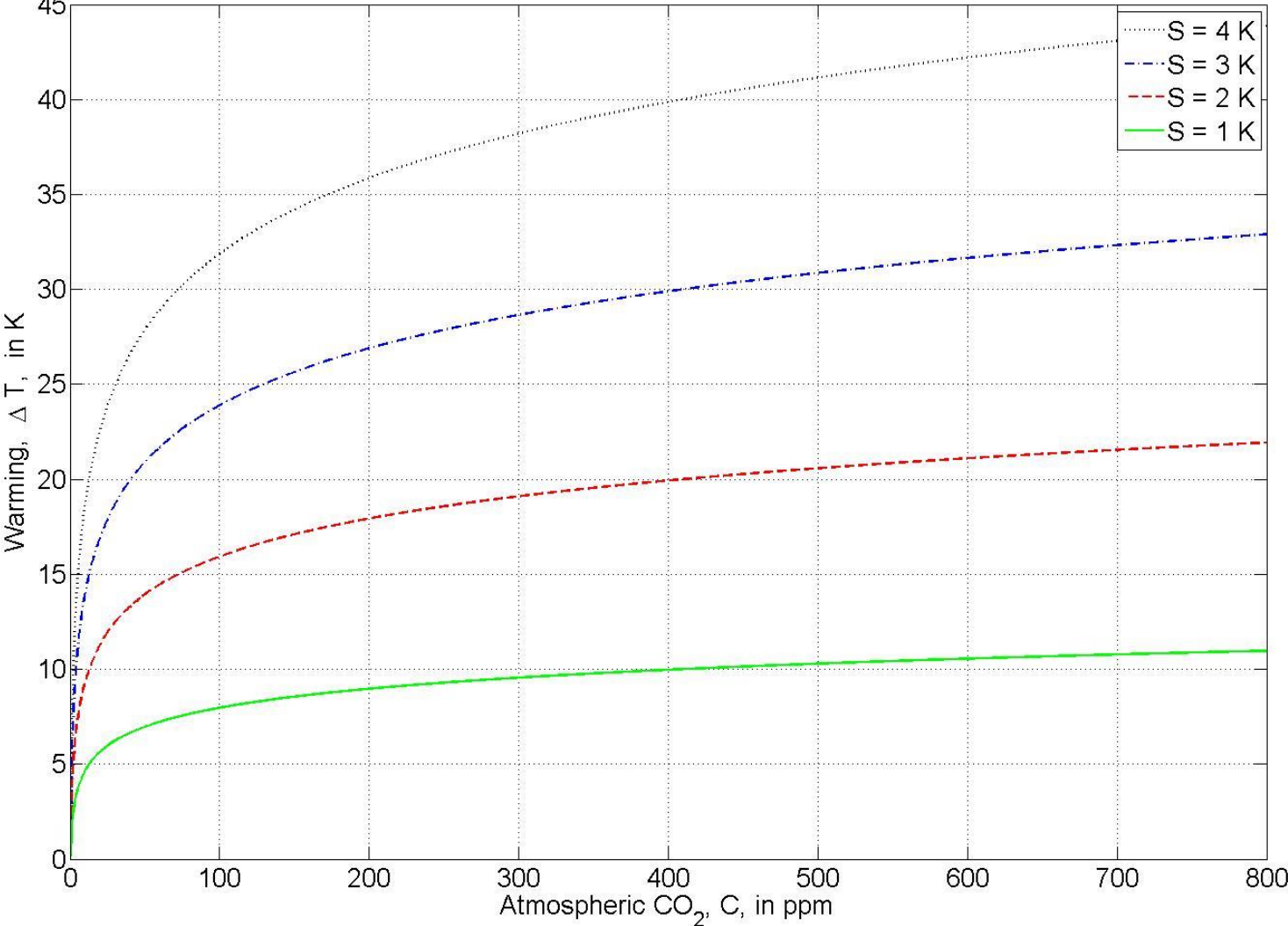


$$\begin{aligned}\Delta T &= \frac{S}{\ln(2)} \ln(C/C_0), \quad \text{for } C > C_0 \\ &= 0, \quad \text{otherwise.}\end{aligned}$$

C = CO₂ concentration; C₀ = saturation concentration; S = Doubling Sensitivity

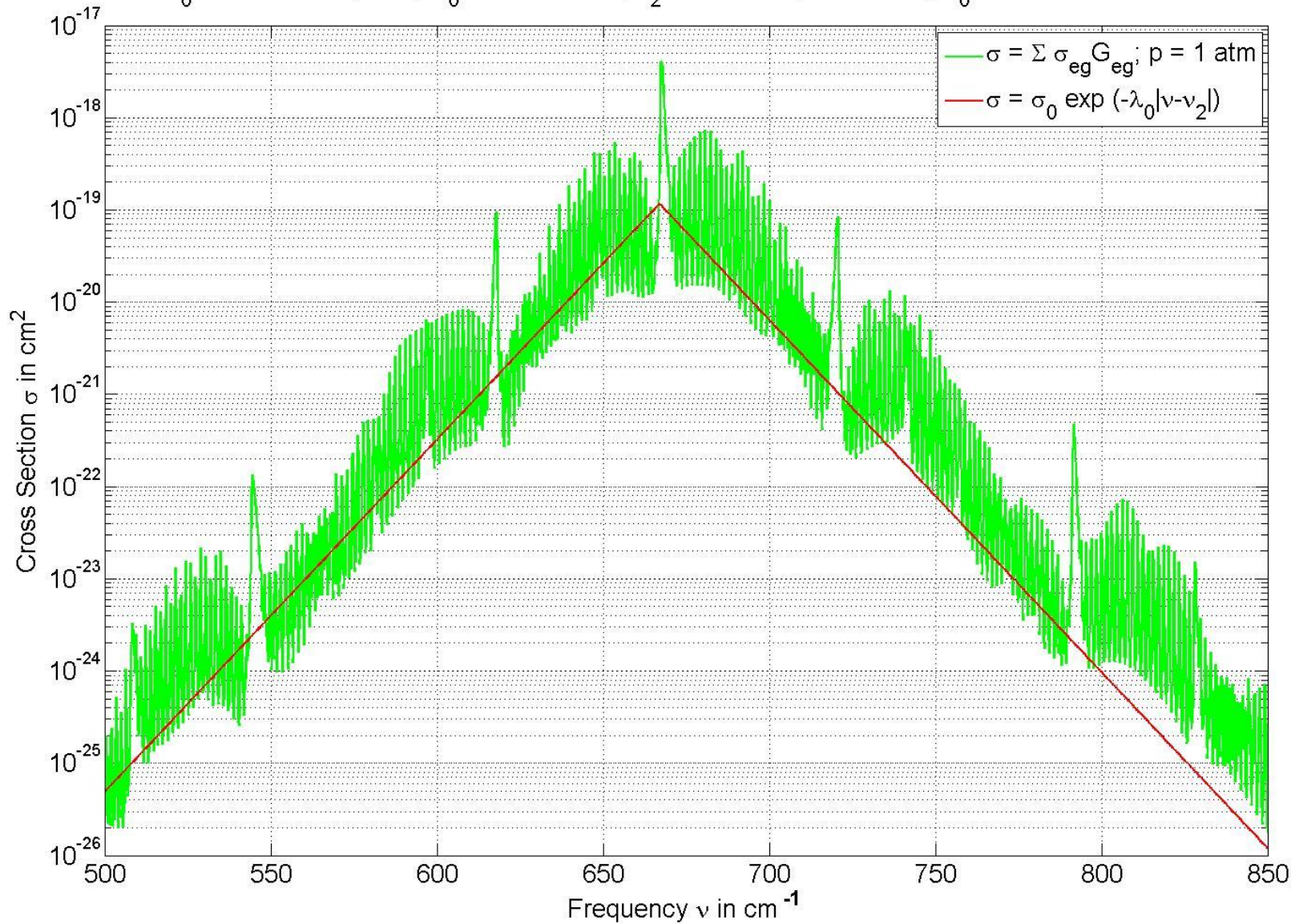
The doubling sensitivity S for greenhouse warming by CO_2 is poorly known. IPCC claims $S = 1.5 \text{ K to } 4.5 \text{ K}$.

$$\Delta T = S \ln(C/C_0)/\ln(2), \text{ if } C > C_0; \quad \Delta T = 0, \text{ otherwise; } \quad C_0 = 0.4 \text{ ppm}$$



Triangular Approximation to CO2 Cross Section Give Logarithmic Warming

$$\sigma_0 = 1.1628e-19; \text{ cm}^2; \lambda_0 = 0.087953 \text{ cm}; \nu_2 = 667 \text{ cm}^{-1}; T = 300 \text{ K}; B_0 = 0.1504 \text{ W m}^{-2} \text{ sr}^{-1}$$



LOGARITHMIC RESPONSE TO CO2 COLUMN DENSITY N

Surface back-radiation from CO2 is very nearly blackbody at surface temperature T_e between upper and lower band limits ν_{\pm} . Multiplying by π for solid angle we find:

$$Q = \pi \int_{\nu_-}^{\nu_+} B(\nu) d\nu \approx \pi(\nu_+ - \nu_-)B_0, \quad \text{with } B_0 = B(\nu_0).$$

Band-edge frequencies ν_{\pm} are where optical depth to space = 1.

$$N\sigma(\nu_{\pm}) = 1 = N\sigma_p e^{-|\nu_{\pm} - \nu_0|/\Delta\nu}.$$

Take logarithms of both sides of equation to find:

$$\nu_+ - \nu_- = 2\Delta\nu \ln N\sigma_p.$$

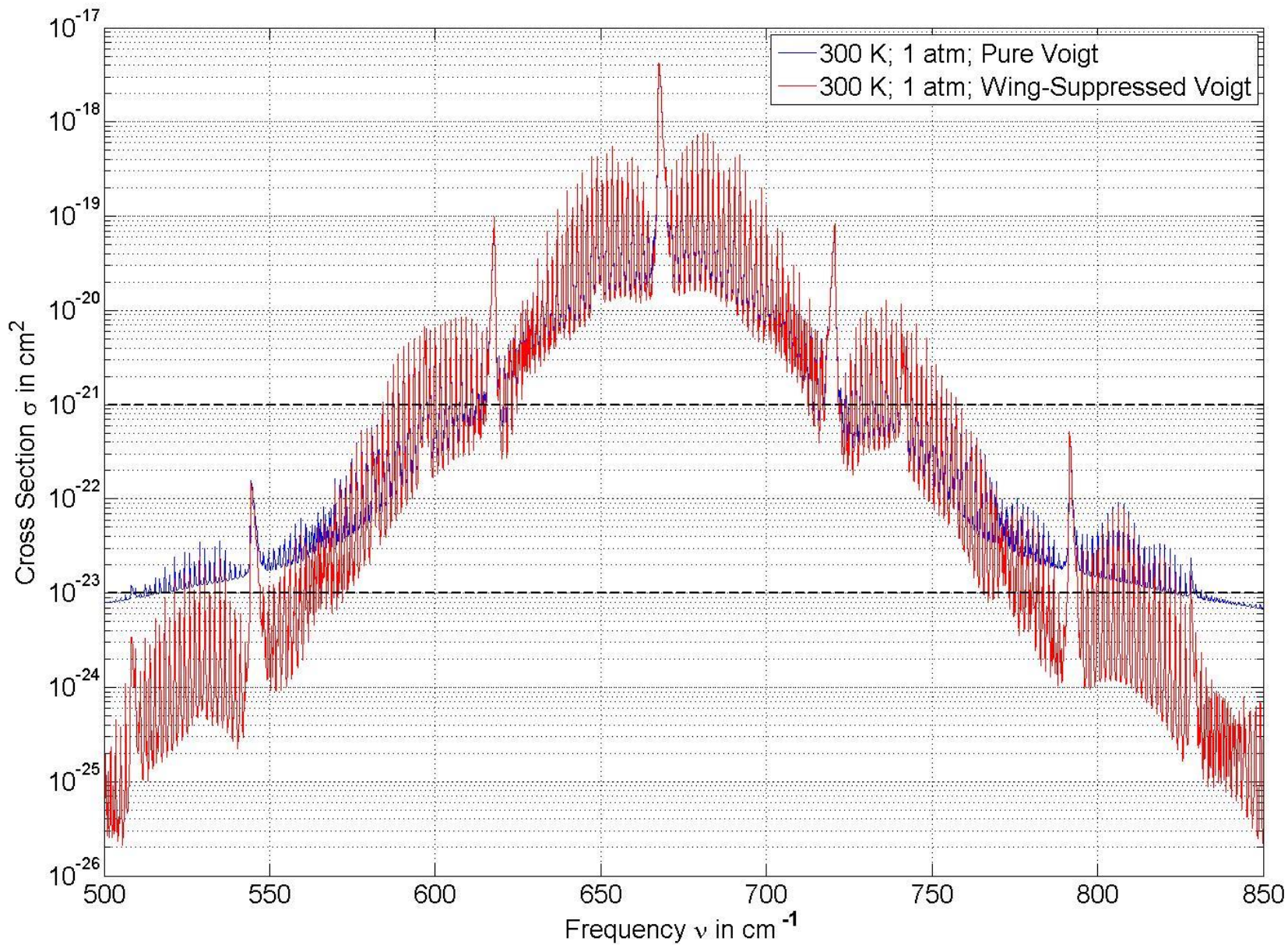
Radiative forcing Q and increment ΔQ for doubling N are:

$$Q = 2\pi B_0 \Delta\nu \ln N\sigma_p, \quad \text{and} \quad \Delta Q = 2\pi B_0 \Delta\nu \ln 2 = 7.4 \text{ W m}^{-2}.$$

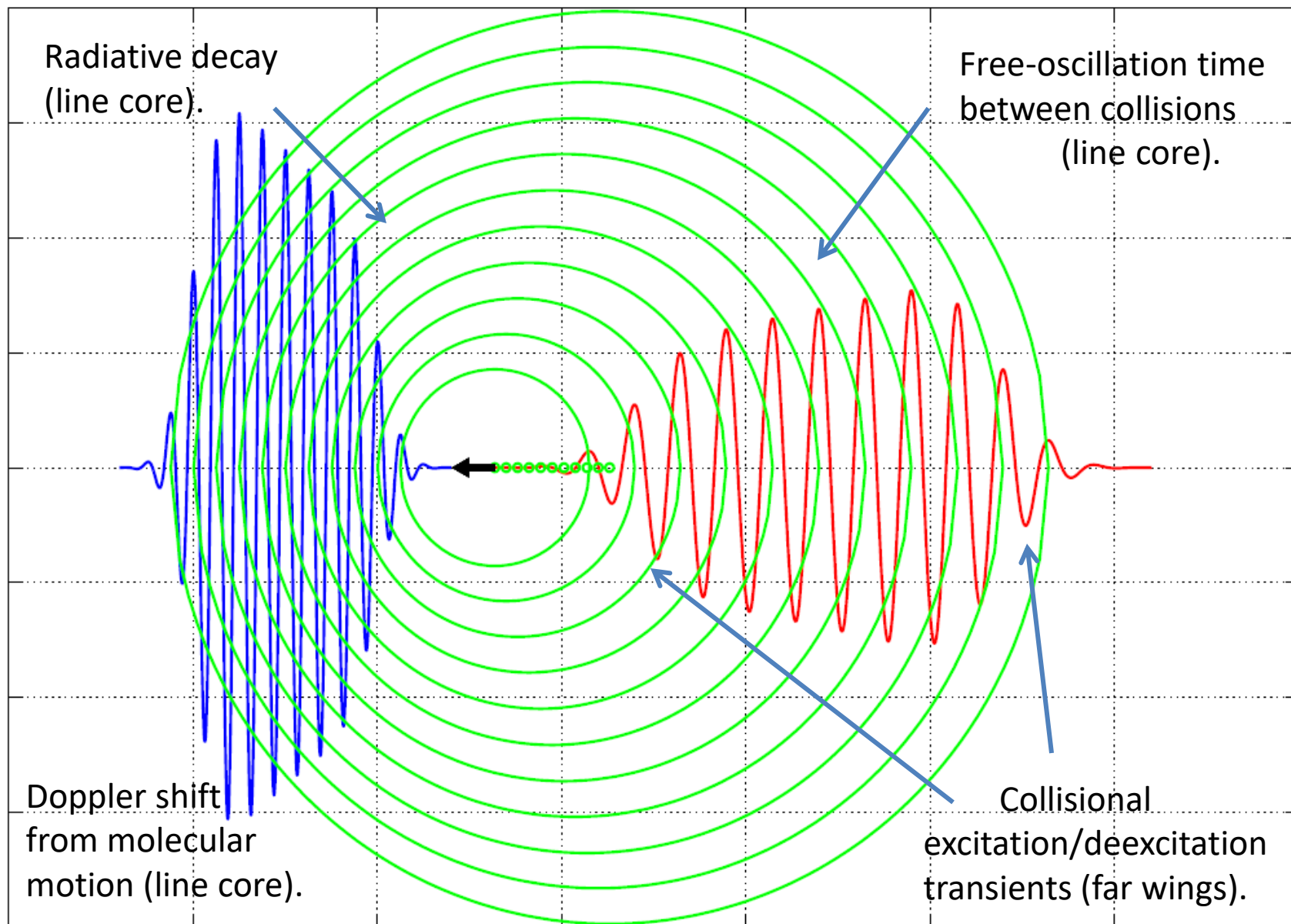
Numbers are $B_0 = 0.14 \text{ W m}^{-2} \text{ str}^{-1}$, $\Delta\nu = 12.2 \text{ cm}^{-1}$. Overlap with the pure-rotational band of water vapor eliminates most of the response from the lower band edge, and IR from clouds further reduces the response to more CO2. The true response is likely less than half the ideal limit or:

$$\Delta Q \leq 3.7 \text{ W m}^{-2}.$$

Cross sections depend on far-wing lineshape at band edges which drive warming!



Physics Behind Shape of Emission Absorption Lines



Radiation

(Pre-QM Fermi's Golden Rule)

Larmor's radiated power

$$P_r = \frac{2q^2 \ddot{X}^2}{3c^3} = m\tau_L \ddot{X}^2$$

The Larmor time

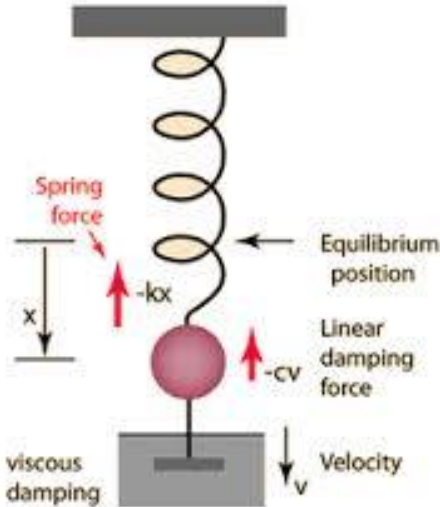
$$\tau_L = \frac{2q^2}{3mc^3} \approx \begin{cases} 10^{-28} \text{ s} & \text{far-infrared,} \\ 10^{-23} \text{ s} & \text{far-ultraviolet.} \end{cases}$$

The Abraham-Lorentz radiation reaction force
(see Jackson's E&M for more discussion)

$$F_r = m\tau_L \ddot{\dot{X}}$$

Newton's second law of motion

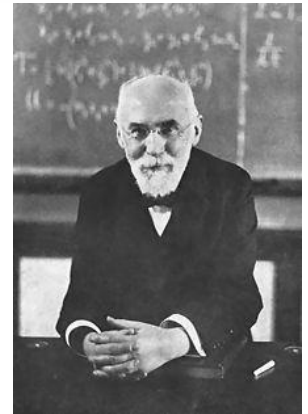
$$m\ddot{X} = -\kappa X + m\tau_L \ddot{\dot{X}} + F$$



Viscous damping does not work as a description of radiation damping. It does not conserve energy. Radiation damping force must be proportional to the third derivative!



Sir Joseph Larmor (1857-1942)



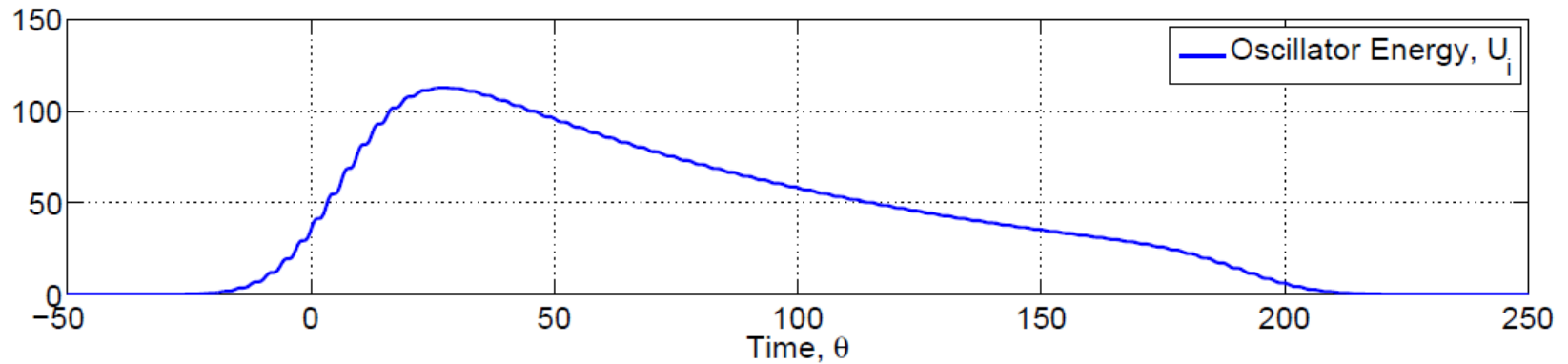
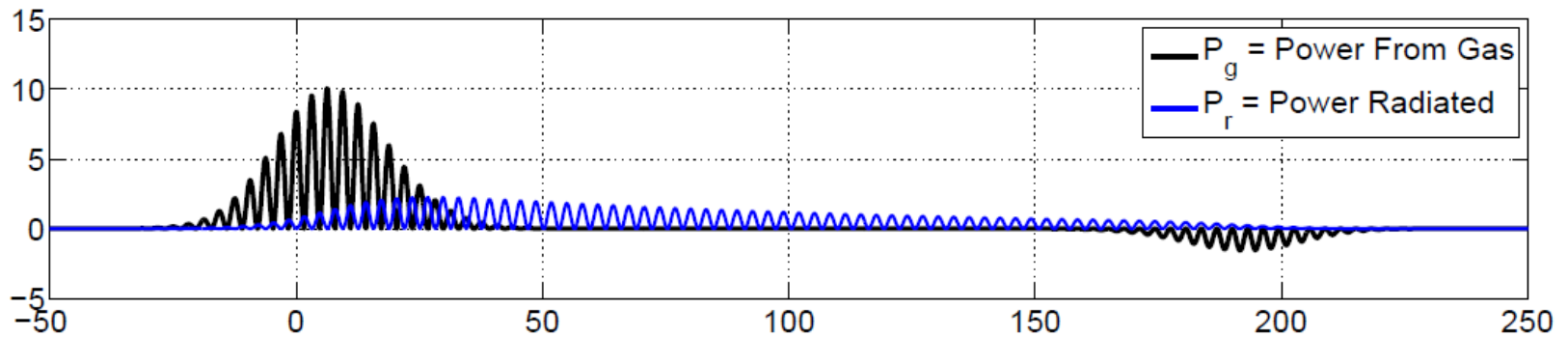
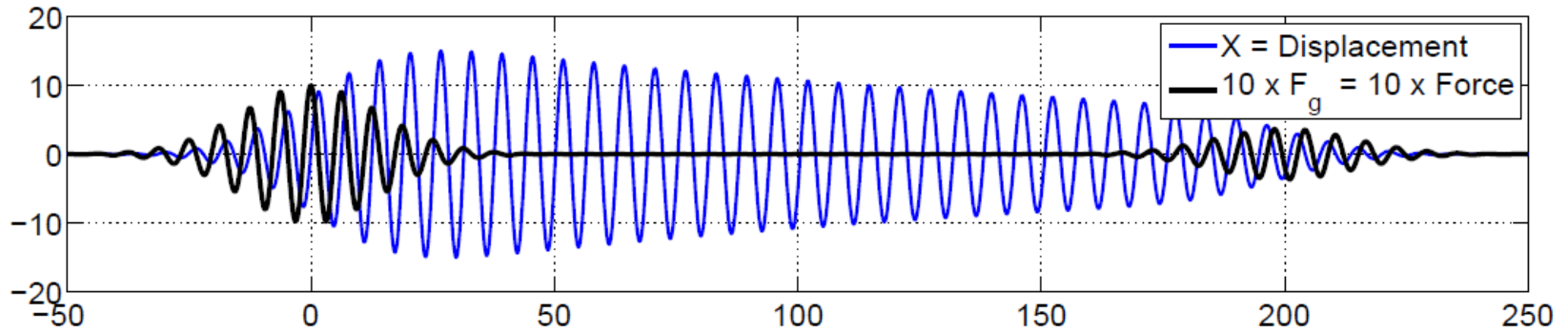
H. A. Lorentz (1853 -1928)



M. Abraham (1875 - 1922)

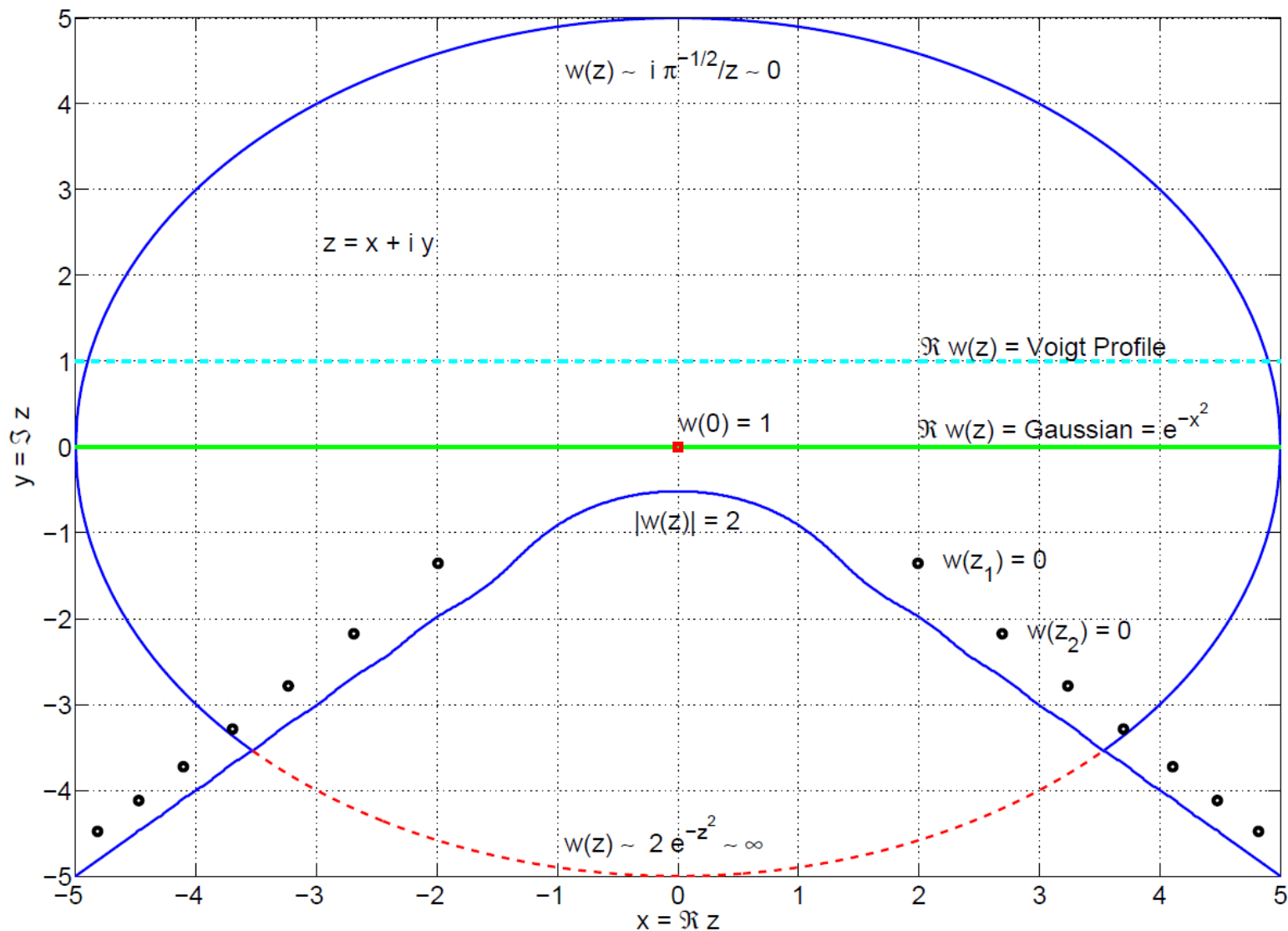
A Solution to Newton's Second Law, $f = ma$

Collisions and Ringing ; $\epsilon_L = 0.01$



A Useful Special Function For Spectral Line Shapes

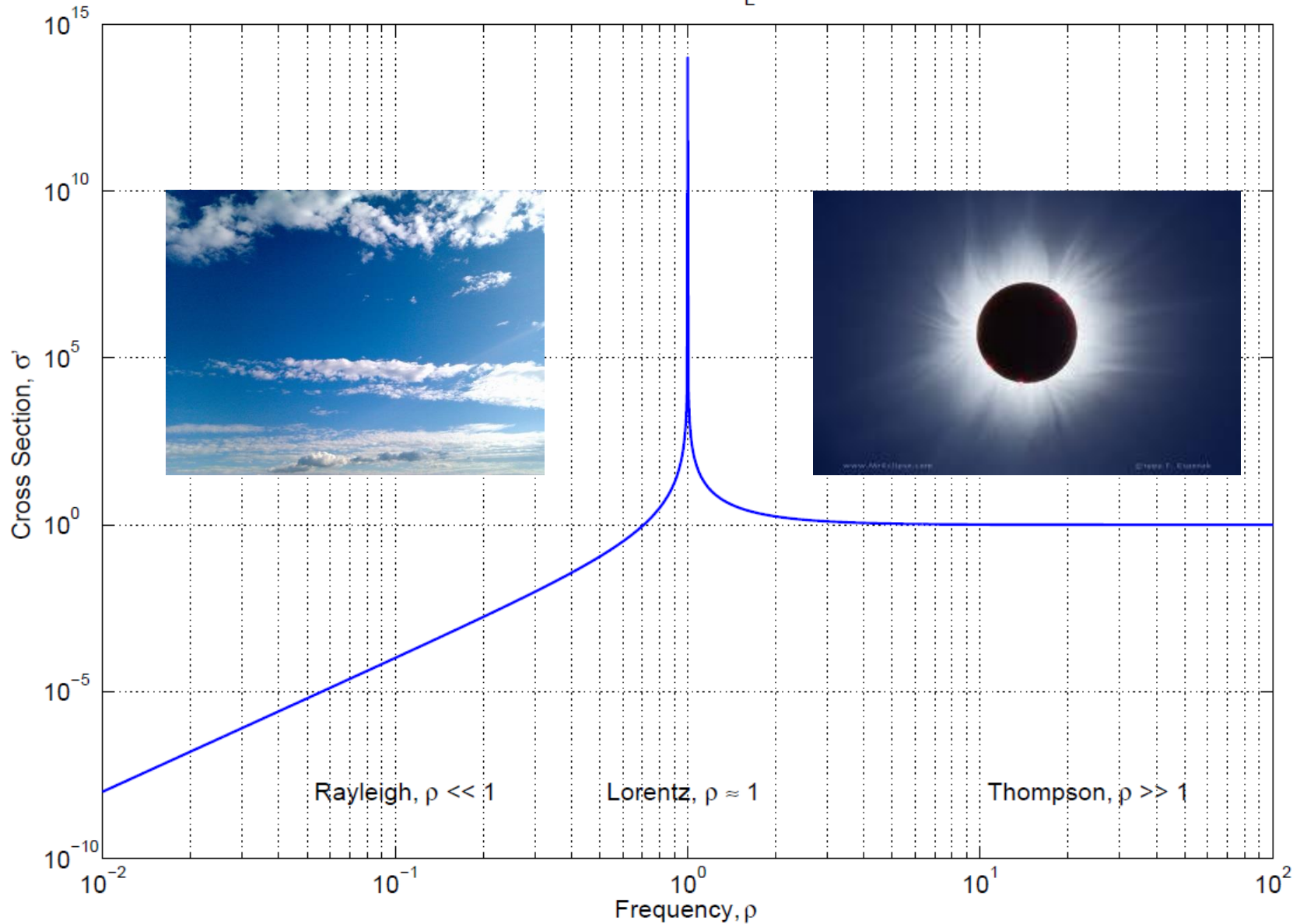
The Faddeeva Function, $w(z) = \pi^{-1/2} \int_0^\infty e^{i z u - u^2/4} du$



Collision-Free Cross Section (not possible with CO₂ in air) :

$$\sigma(\rho) = \frac{\rho^4}{(1 - \rho^2)^2 + \epsilon_L^2 \rho^6}$$

Cross Sections for $\epsilon_L = 1e-07$



Outline of Lineshape Calculation with Collisions

$$\begin{aligned}\text{Relative time:} & \quad \theta = \omega_0 t \\ \text{Relative frequency:} & \quad \rho = \omega/\omega_0\end{aligned}$$

$$\begin{aligned}\text{Collisional Force } F &= 2 \cos(\rho_k \theta) A(\theta) \\ &= \text{carrier of frequency } \rho_k \times \text{envelope } A(\theta)\end{aligned}$$

The Fourier transform of the envelope is

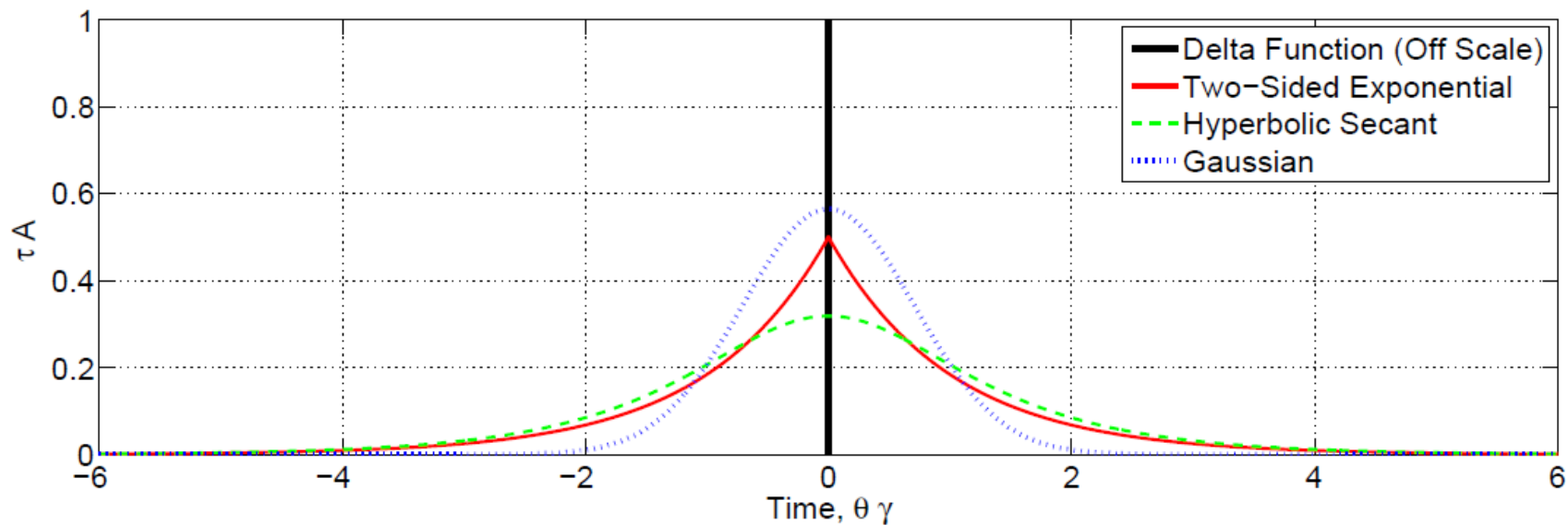
$$\tilde{A}(\rho) = \int_{-\infty}^{\infty} A(\theta) e^{i\rho\theta} d\theta,$$

The resulting lineshape function is a Lorentzian, weighted by the absolute squared Fourier transform of the envelope:

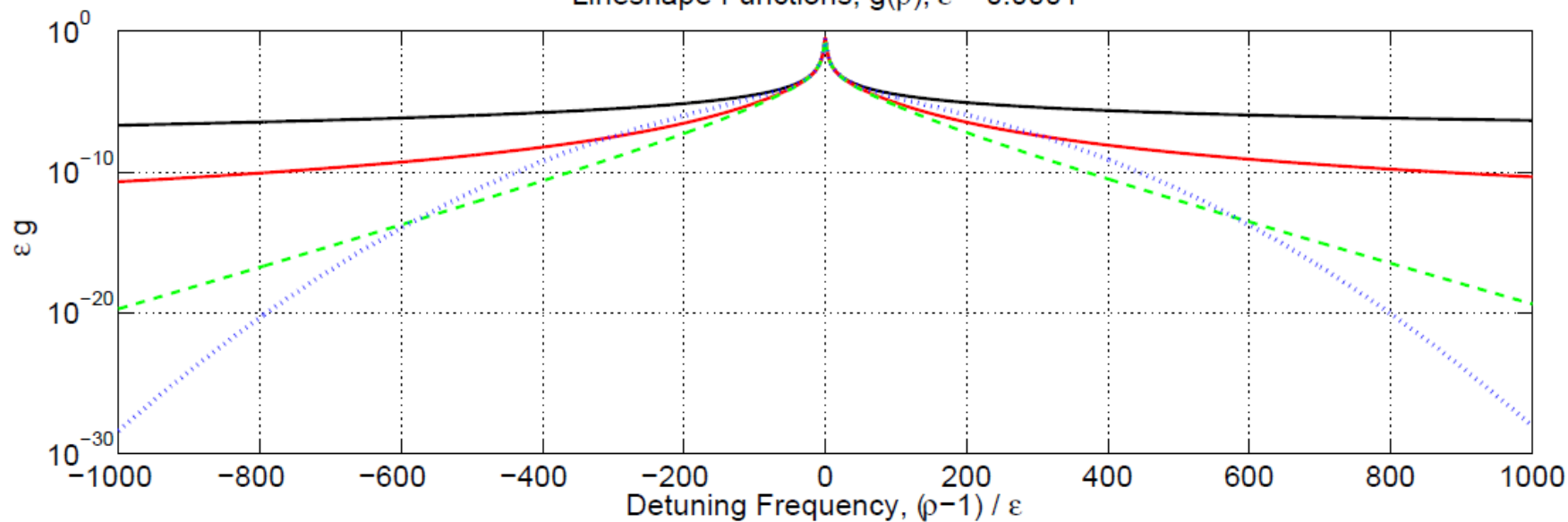
$$g(\rho) = \frac{\epsilon \rho^4 |\tilde{A}(\rho - 1)|^2}{\pi \{(\rho - 1)^2 + \epsilon^2\}}.$$

$\epsilon = \text{natural radiative broadening} + \text{collisions}$

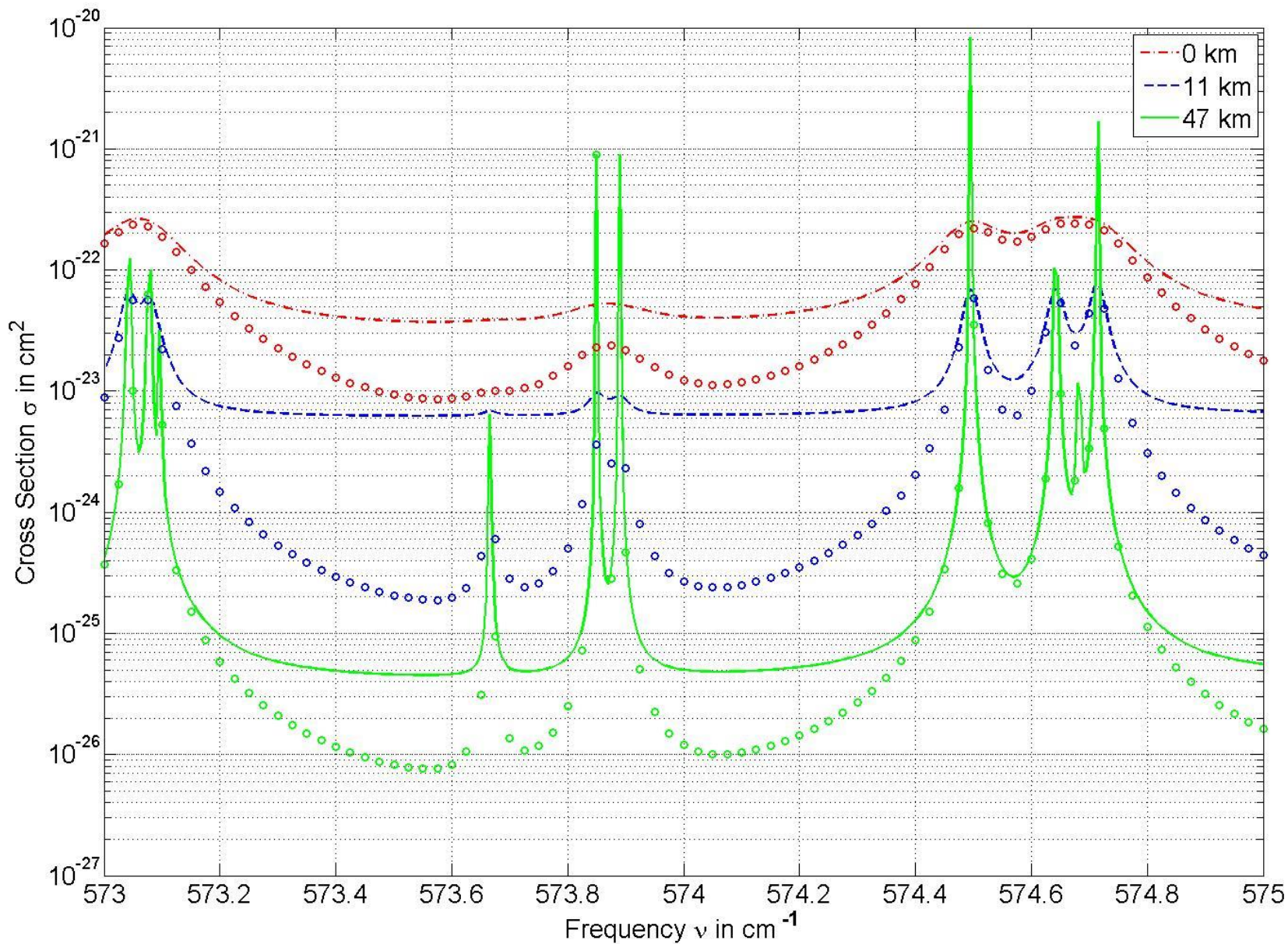
Collisional Force Envelopes, $A(\theta)$; $\tau = 1/\gamma = 100$



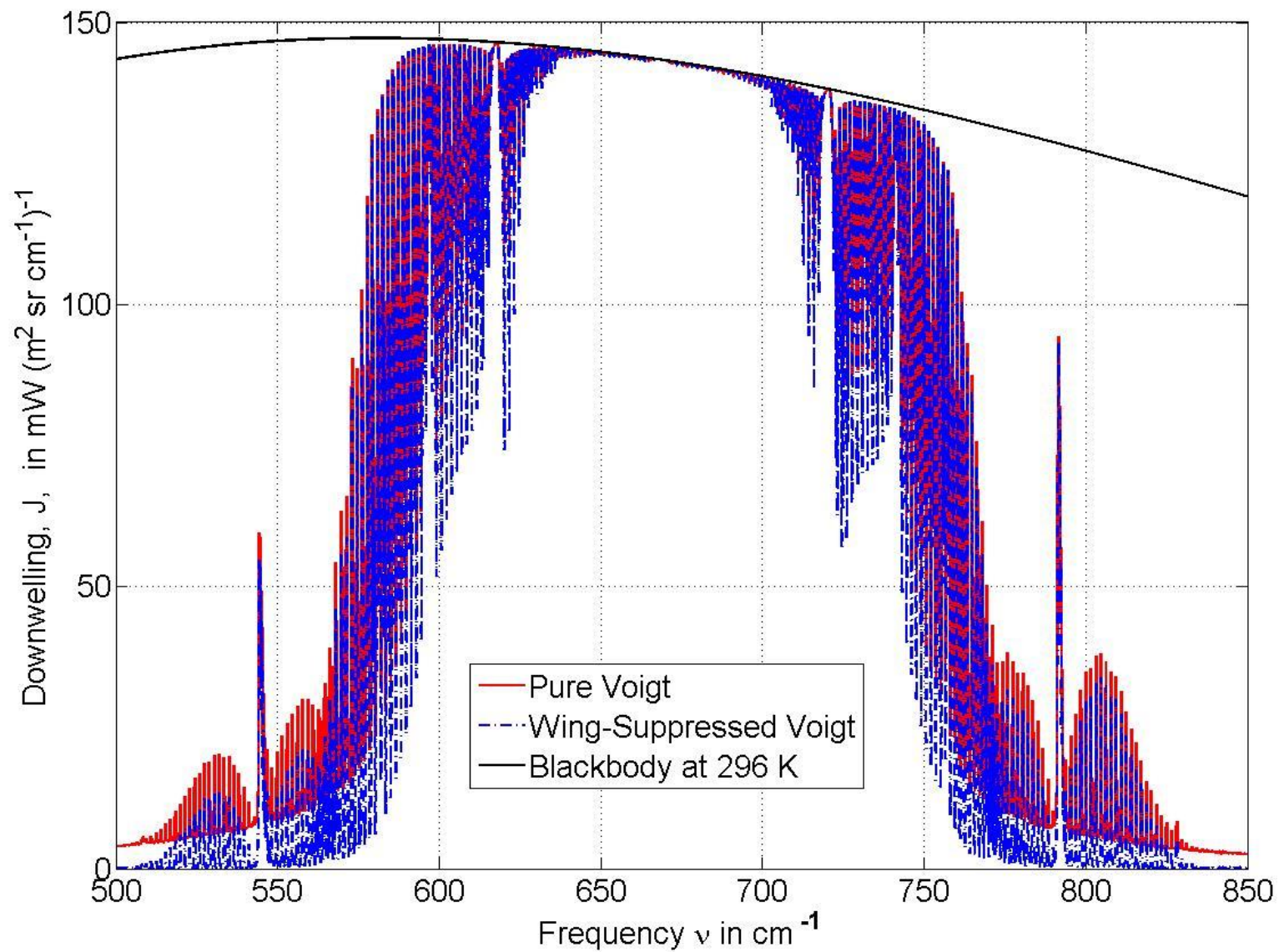
Lineshape Functions, $g(\rho)$; $\varepsilon = 0.0001$



Far-wing (global warming) cross sections much bigger with Lorentz broadening (lines) than with realistic far-wing broadening (circles).



390 ppm CO₂



Using Voigt profiles increases
the radiative-forcing increment
from doubling CO₂ by a factor
~1.4

But far wing absorption from Voigt
profiles does not exist!

Need experimental measurements!

Voigt Line Shapes Don't Work in Far Wings!

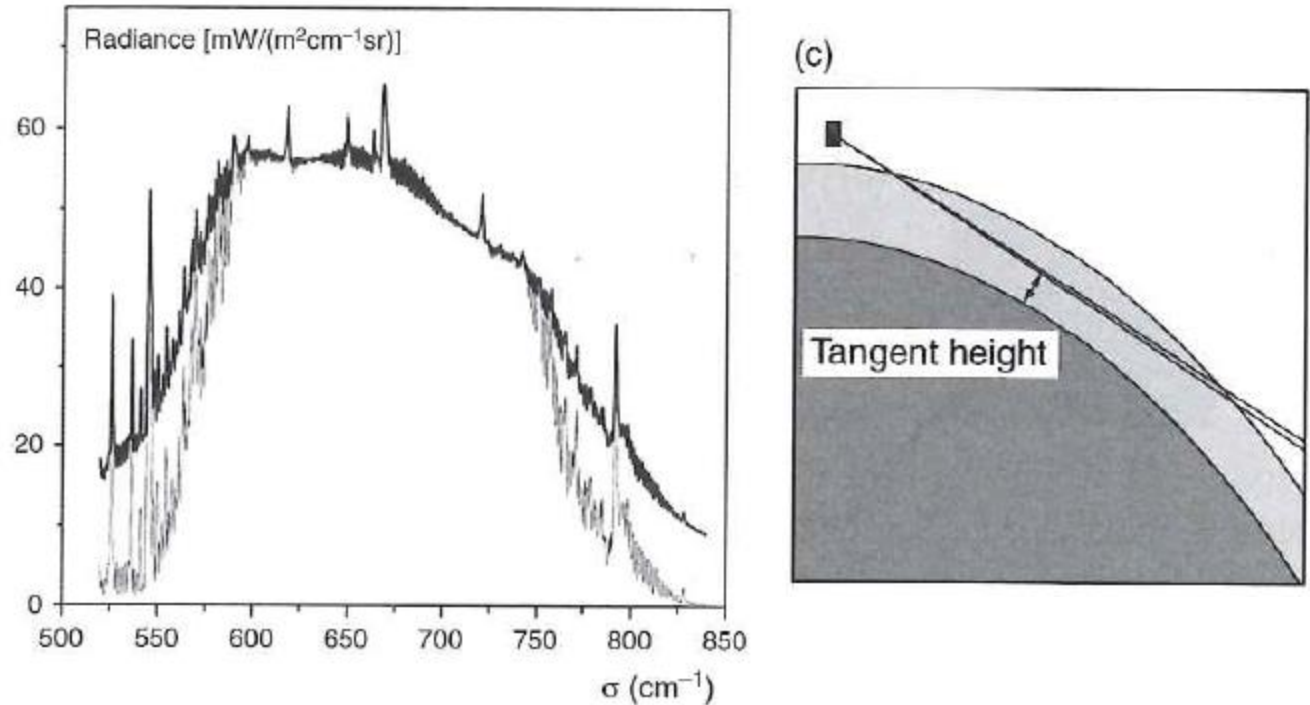


Fig. VII.15: Radiance emitted by the Earth atmosphere in the region of the ν_2 CO₂ band for a 1 cm^{-1} resolution. The thin line gives measured values obtained by a balloon-borne instrument²⁹ at 40 km altitude looking down to a tangent height of 10 km. The thick line corresponds to predictions using purely Voigt line shapes. After Ref. 603.



Climate models are not working! Far-wing lineshapes (forcing ΔQ_2) are one of many possible causes. Clouds (feedback f) probably even more important

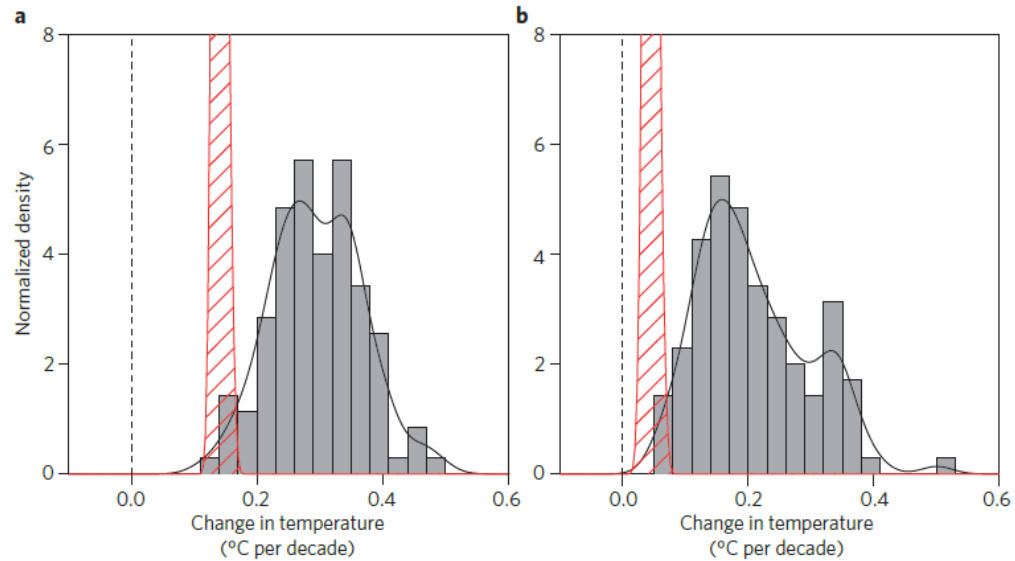


Figure 1 | Trends in global mean surface temperature. **a**, 1993–2012. **b**, 1998–2012. Histograms of

“Science is the belief in the ignorance of experts.”

Richard Feynman